

## THE $\eta(1405)$ , $\eta(1475)$ $f_1(1420)$ , AND $f_1(1510)$

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The first observation of the  $\eta(1440)$  was made in  $p\bar{p}$  annihilation at rest into  $\eta(1440)\pi^+\pi^-$ ,  $\eta(1440) \rightarrow K\bar{K}\pi$  (BAILLON 67). This state was reported to decay through  $a_0(980)\pi$  and  $K^*(892)\bar{K}$  with roughly equal contributions. The  $\eta(1440)$  was also observed in radiative  $J/\psi(1S)$  decay to  $K\bar{K}\pi$  (SCHARRE 80, EDWARDS 82E, AUGUSTIN 90). There is now evidence for the existence of two pseudoscalars in this mass region, the  $\eta(1405)$  and  $\eta(1475)$ . The former decays mainly through  $a_0(980)\pi$  (or direct  $K\bar{K}\pi$ ) and the latter mainly to  $K^*(892)\bar{K}$ .

The simultaneous observation of two pseudoscalars is reported in three production mechanisms:  $\pi^-p$  (RATH 89, ADAMS 01); radiative  $J/\psi(1S)$  decay (BAI 90C, AUGUSTIN 92);  $\bar{p}p$  annihilation at rest (BERTIN 95, BERTIN 97, CICALO 99, NICHITIU 02). All of them give values for the masses, widths and decay modes in reasonable agreement. However, AUGUSTIN 92 favors a state decaying into  $K^*(892)\bar{K}$  at a lower mass than the state decaying into  $a_0(980)\pi$ , although agreement with MARK-III is not excluded. In  $J/\psi(1S)$  radiative decay, the  $\eta(1405)$  decays into  $K\bar{K}\pi$  through  $a_0(980)\pi$  and hence a signal is also expected in the  $\eta\pi\pi$  mass spectrum. This was indeed observed by MARK III in  $\eta\pi^+\pi^-$  (BOLTON 92B) which report a mass of 1400 MeV, in line with the existence of the  $\eta(1405)$  decaying to  $a_0(980)\pi$ . This state is also observed in  $\bar{p}p$  annihilation at rest into  $\eta\pi^+\pi^-\pi^0\pi^0$ , where it decays into  $\eta\pi\pi$  (AMSLER 95F). The intermediate  $a_0(980)\pi$  accounts for roughly half of the  $\eta\pi\pi$  signal, in agreement with MARK III (BOLTON 92B) and DM2 (AUGUSTIN 90).

The existence of the  $\eta(1295)$  is questioned by KLEMP 05. However, this state has been observed by four  $\pi^-p$  experiments (ADAMS 01, FUKUI 91C, ALDE 97B, MANAK 00A) and evidence is also reported in  $\bar{p}p$  annihilation (ABELE 98, ANISOVICH 01, AMSLER 04B). In  $J/\psi$  radiative decay an  $\eta(1295)$  signal is seen in the  $0^{-+}$   $\eta\pi\pi$  wave of DM2 data (AUGUSTIN 92).

Assuming that the  $\eta(1295)$  is established, the  $\eta(1475)$  could be the first radial excitation of the  $\eta'$ , with the  $\eta(1295)$  being the first radial excitation of the  $\eta$ . Ideal mixing, suggested by the  $\eta(1295)$  and  $\pi(1300)$  mass degeneracy, would then imply that the second isoscalar in the nonet is mainly  $s\bar{s}$  and hence couples to  $K^*\bar{K}$ , in agreement with the  $\eta(1475)$ . Also its width matches the expected width for the radially excited  $s\bar{s}$  state (CLOSE 97, BARNES 97).

The  $K\bar{K}\pi$  and  $\eta\pi\pi$  channels were studied in  $\gamma\gamma$  collisions (ACCIARRI 01G). The analysis leads to an  $\eta(1475)$  signal in  $K\bar{K}\pi$ , but the  $\eta(1405)$  is not observed in  $K\bar{K}\pi$  nor in  $\eta\pi\pi$ . This result is somewhat in disagreement with CLEO-II which did not observe any pseudoscalar signal in  $\gamma\gamma \rightarrow \eta(1475) \rightarrow K_S^0 K^\pm \pi^\mp$  (AHOHE 05), but more data are required. Since gluonium production is presumably suppressed in  $\gamma\gamma$  collisions, the ACCIARRI 01G results suggest that the  $\eta(1405)$  has a large gluonic content (see also CLOSE 97B, LI 03C). The observation of the  $\eta(1475)$  combined with the absence of an  $\eta(1405)$  signal strengthens the two-resonances hypothesis.

The gluonium interpretation is not favored by lattice gauge theories which predict the  $0^{-+}$  state above 2 GeV (BALI 93). However, the  $\eta(1405)$  is an excellent candidate for the  $0^{-+}$  glueball in the flux tube model (FADDEEV 04). In this model the  $0^{++}$   $f_0(1500)$  glueball is also naturally related to a  $0^{-+}$  glueball with mass degeneracy broken in QCD.

Let us now deal with  $1^{++}$  isoscalars. The  $f_1(1420)$ , decaying to  $K^*\bar{K}$ , was first reported in  $\pi^-p$  reactions at 4 GeV/c (DIONISI 80). However, later analyses found that the 1400–1500 MeV region was far more complex (CHUNG 85, REEVES 86, BIRMAN 88). A reanalysis of the MARK III data in radiative  $J/\psi(1S)$  decay to  $K\bar{K}\pi$  (BAI 90C) shows the  $f_1(1420)$  decaying into  $K^*\bar{K}$ . Also, a  $C=+1$  state is observed in tagged  $\gamma\gamma$  collisions (*e.g.*, BEHREND 89).

In  $\pi^-p \rightarrow \eta\pi\pi n$  charge-exchange reactions at 8–9 GeV/c the  $\eta\pi\pi$  mass spectrum is dominated by the  $\eta(1295)$  and  $\eta(1440)$  (ANDO 86, FUKUI 91C) and at 100 GeV/c ALDE 97B report the  $\eta(1295)$  and  $\eta(1440)$  decaying to  $\eta\pi^0\pi^0$ , with a weak  $f_1(1285)$  signal and no evidence for the  $f_1(1420)$ .

Axial ( $1^{++}$ ) mesons are not observed in  $\bar{p}p$  annihilation at rest in liquid hydrogen, which proceeds dominantly through  $S$ -wave annihilation. However, in gaseous hydrogen  $P$ -wave annihilation is enhanced and, indeed, BERTIN 97 report  $f_1(1420)$  decaying to  $K^*\bar{K}$ .

The  $f_1(1420)$ , decaying into  $K\bar{K}\pi$ , is also seen in  $pp$  central production together with the  $f_1(1285)$ . The latter decays via  $a_0(980)\pi$  and the former only via  $K^*\bar{K}$ , while the  $\eta(1440)$  is absent (ARMSTRONG 89, BARBERIS 97C). The  $K_S K_S \pi^0$  decay mode of the  $f_1(1420)$  establishes unambiguously  $C=+1$ . On the other hand, there is no evidence for any state decaying to  $\eta\pi\pi$  around 1400 MeV and hence the  $\eta\pi\pi$  mode of  $f_1(1420)$  must be suppressed (ARMSTRONG 91B).

We now turn to the experimental evidence for the  $f_1(1510)$ . Two states, the  $f_1(1420)$  and the  $f_1(1510)$ , decaying to  $K^*\bar{K}$ , compete for the  $s\bar{s}$  assignment in the  $1^{++}$  nonet. The  $f_1(1510)$  was seen in  $K^-p \rightarrow \Lambda K\bar{K}\pi$  at 4 GeV/ $c$  (GAVILLET 82) and at 11 GeV/ $c$  (ASTON 88C). Evidence is also reported in  $\pi^-p$  at 8 GeV/ $c$ , based on the phase motion of the  $1^{++}$   $K^*\bar{K}$  wave (BIRMAN 88).

The absence of  $f_1(1420)$  in  $K^-p$  (ASTON 88C) argues against this state being the  $s\bar{s}$  member of the  $1^{++}$  nonet. However, the  $f_1(1420)$  was reported in  $K^-p$  but not in  $\pi^-p$  (BITYUKOV 84) while two experiments do not observe the  $f_1(1510)$  in  $K^-p$  (BITYUKOV 84, KING 91). It is also not seen in radiative  $J/\psi(1S)$  decay (BAI 90C, AUGUSTIN 92), central collisions (BARBERIS 97C), nor in  $\gamma\gamma$  collisions (AIHARA 88C), although, surprisingly for an  $s\bar{s}$  state, a signal is reported in  $4\pi$  decays (BAUER 93B). These facts lead to the conclusion that the  $f_1(1510)$  needs experimental confirmation (CLOSE 97D).

Assigning the  $f_1(1420)$  to the  $1^{++}$  nonet one finds a nonet mixing angle of  $\sim 50^\circ$  (CLOSE 97D). However, arguments favoring the  $f_1(1420)$  being a hybrid  $q\bar{q}g$  meson or a four-quark state were put forward by ISHIDA 89 and by CALDWELL 90, respectively, while LONGACRE 90 argued for a molecular state formed by the  $\pi$  orbiting in a  $P$ -wave around an  $S$ -wave  $K\bar{K}$  state.

Summarizing, there is convincing evidence for the  $f_1(1420)$  decaying to  $K^*\overline{K}$ , and for two pseudoscalars in the  $\eta(1440)$  region, the  $\eta(1405)$  and  $\eta(1475)$ , decaying to  $a_0(980)\pi$  and  $K^*\overline{K}$ , respectively. The  $f_1(1510)$  is not well established.