# DATA ON ELEMENTARY PARTICLES AND RESONANT STATES, NOVEMBER 1963 

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#### Abstract

The elementary particle data comprise quantum numbers, mass, magnetic moment, mean life, common decay modes and branching ratios. The resonant state data comprise quantum numbers, mass, width, life time, threshold laboratory momentum for a production process, decay modes, branching ratios and $Q$ values. All data used to obtain the tabulated average values, and some data not used are given in 52 footnotes.


## 1. Introduction

This is a revised version of previously published tables (Roos, 1963), dated March 1963. In the previous version no distinction was made between generally accepted resonances and not (yet) accepted or dubious resonances.

In this respect the present tables are radically different. First we have omitted all such resonances, which were considered dubious a year ago, and for which no further evidence has been produced since then. This selection rule excludes the resonances named $\mathrm{Y}^{*}, \mathrm{Y}_{0}^{* *}, \mathrm{Z}_{3}^{*}, \mathrm{~N}_{11}^{*}, \mathrm{Z}_{1}^{*}, \mathrm{~K}_{3}^{*}, \chi_{2}, \kappa_{3}, \mathrm{~K}^{* *}, \mathrm{~K}_{5}^{*}, \chi_{1}, \psi_{5}, \phi_{3}, \rho_{2}, \rho_{1}, \psi_{4}, \delta, \alpha$, $\psi_{3}, \zeta, \psi_{2}, \phi_{1}$ and $\psi_{1}$ in the previous version. Some of the evidence for the above resonances is still used as possible evidence for particular decay modes of accepted resonances.

Further we have attempted to divide the accepted resonances into two confidence classes. The higher confidence class contains the generally accepted resonances, while the lower class contains dubious resonances or very new resonances for which the data still are meager. The difference in notation between the two classes is very simple: the accepted resonances have symbols, while the lower confidence class is left without symbols. From the previous tables the resonances $\mathrm{Y}^{* *}, \mathrm{Y}_{2}^{*}, \kappa_{2}$ and $\phi_{2}$ have been stripped of their symbols and shuffled into the lower class.

Table 1 contains the particles which are customarily called elementary, and their antiparticles.

Table 2 contains the baryonic resonant states, defined as states with baryonic number $B=1$. The generic symbols $\Xi^{*}, \mathrm{Y}^{*}$ and $\mathrm{N}^{*}$ are used for states with strangeness $-2,-1$ and 0 , respectively, and they carry two subscripts, corresponding to isospin and spin in the following manner:

$$
\Xi_{2 T, 2 J}^{*}, \mathrm{Y}_{T, 2 J}^{*}, \mathrm{~N}_{2 T, 2 J}^{*}
$$

Not known subscripts are left blank. Further degeneracy is resolved by adding more stars to the heavier resonances. Table 3 contains the mesonic resonant states, defined as states with baryonic number $B=0$.

## 2. Quantum Numbers

A blank space in any of the quantum-number columns may signify that the quantity in question is not known ${ }^{\dagger}$ or that it cannot be defined ( $T, T_{3}, S$, and parity for leptons). $T$ is not repeated for isospin multiplets, nor is $T, S$, or parity repeated for antiparticles. In Tables 2 and 3, $J$, parity and $G$ parity are not repeated for different charge states.

Different charge states are given separate entries when they have been found.
Parity is defined in relation to N and K ; by definition, N has parity + and K has-.

## 3. Mass and Magnetic Moment

The mass is given in two units, MeV and $m_{\pi^{ \pm}}$; for leptons $m_{\pi^{ \pm}}$is exchanged for $m_{\mathrm{e}}$ We choose the units

$$
\begin{aligned}
\hbar & =c=1 \\
m_{\pi^{ \pm}} & =139.58 \mathrm{MeV}=2.4881 \times 10^{-25} \mathrm{~g} \\
m_{\mathrm{e}} & =0.510976 \mathrm{MeV}
\end{aligned}
$$

The accurate mass is expressed in MeV for all particles, except the muon and the neutrinos, for which it is expressed in $m_{\mathrm{e}}$. The mass in the other unit gives only the significant figures.

The magnetic moment is expressed in proton magnetons for baryons, in muon magnetons for the muon, and in electron magnetons for the electron.

The mass and magnetic moments of antiparticles are included when they have been specifically measured; otherwise, a blank space is left. A blank space in the magneticmoment column may also indicate that the value is not known.

## 4. Width and Lifetime

In table 1 the mean life is given in two units, accurately in seconds and rounded off in $1 / m_{\pi^{ \pm}}$. The relation is

$$
1 / m_{\pi^{ \pm}}=4.7153 \times 10^{-24} \mathrm{~s}
$$

the latter time signifying the time required for light to travel the distance of a Compton wavelength of the $\pi^{ \pm}$-meson. This distance equals

$$
1.4136 \times 10^{-13} \mathrm{~cm}
$$

In tables 2 and 3 the full width $\Gamma$ at half-maximum of the resonance is given in

[^0]MeV , and the lifetime $\Gamma^{-1}$ in units of $1 / \mathrm{m}_{\pi^{\star}}$, to allow comparison with the mean lives in table 1 .

A blank space means that the quantity is not known. Widths and lifetimes of antiparticles, and, in tables 2 and 3 charge multiplets, are not included unless they have been specifically measured.

## 5. Production Properties

In tables 2 and 3 one production reaction is given although others may also have been used. The laboratory momentum of the incident particle has been computed for that production reaction, at the threshold of resonance production.

A blank space in the $k_{\text {lab }}$ column signifies that the rest masses in the production reaction have not been computed because no threshold exists or because it is questionable which threshold is of interest.

A blank space in both production columns indicates that detailed information on the production of different charge states is not available.

## 6. Decay Properties

The commonest decay modes are given if they have been observed. By "commonest" we mean a branching ratio $\geqq 1 \%$.

The decay modes of antiparticles are not listed because they are simply the antimodes of the particles, and the branching ratios are the same. A blank space in any of the decay-property columns signifies that the information is lacking.

## 7. Footnotes and References

All detailed information is collected in the footnotes, where, however, references are given only in the contracted form: first author (year). The full references are collected alphabetically in a separate list. Use has been made of all literature available in Copenhagen by November 1963, and of all results presented in the strong interaction parallel sessions or in the plenary sessions at the Sienna International Conference on Elementary Particles, September 30. - October 5., Sienna, Italy (1963).

The compiler is indebted to countless people for the privilege of being informed about experimental results prior to publication. It is a pleasure to acknowledge the help given by members of the staffs of NORDITA and the Institute for Theoretical Physics at the University of Copenhagen, as well as the financial support of NORDITA.
Table 1
Elementary particles,


Z әТ甲ец

|  |  | $\stackrel{\sim}{\sim}$ | \＃ | N | ๙ | H | $\stackrel{\leftrightarrow}{\circ}$ | $\stackrel{\circ}{\sim}$ | N |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\begin{gathered} \text { む̀ } \\ \stackrel{\circ}{\circ} \end{gathered}$ | © | © | $\underset{\infty}{\infty}$ | $\underset{\sim}{\#} \stackrel{n}{F}$ | ొ | 込 | 우 | か～が |  |
|  |  |  |  |  |  |  | － |  |  |
|  | \％ | $\begin{aligned} & 1 \\ & + \\ & + \\ & + \\ & + \\ & \hline \end{aligned}$ |  |  |  |  | 旨 |  |  |
|  |  | 俞 | 苍 | 各 | 雨 | $\stackrel{1}{\mathrm{~N}}$ | $\stackrel{\circ}{\text { N }}$ | $\stackrel{\text { N }}{\stackrel{\sim}{n}}$ | セ8 |
|  | $\begin{aligned} & \text { O. } \\ & \text { 品 } \\ & 0 \end{aligned}$ | $\underset{\sim}{\infty}$ | $\begin{aligned} & 2 \\ & 4 \end{aligned}$ | $\begin{gathered} z \\ z \\ i \end{gathered}$ | $\begin{gathered} z \\ b \\ i \end{gathered}$ |  | $\begin{aligned} & n_{1} \\ & k \\ & k \end{aligned}$ |  | 4 |
|  |  |  | N | $\stackrel{\Gamma}{i}$ | $\begin{aligned} & \infty \\ & \text { in } \end{aligned}$ | $\stackrel{\sim}{\infty}$ | $\vec{i}$ | 운 | $\stackrel{\sim}{\infty}$ |
|  |  |  | 은 | $\stackrel{\substack{\infty \\ \mathrm{V}}}{ }$ | 8 | $\begin{aligned} & 0 \\ & 0 \\ & +1 \\ & 0 \\ & \hline \end{aligned}$ | $\underset{\sim}{\underset{\sim}{\sim}}$ | N + $\sim$ | N + + $\bullet$ $\bullet$ |
|  | $\begin{aligned} & \text { F } \\ & \text { E } \end{aligned}$ | $\begin{aligned} & \infty \\ & \stackrel{\oplus}{-1} \end{aligned}$ | $\begin{aligned} & 0 \\ & \underset{\sim}{0} \end{aligned}$ | $\begin{aligned} & \stackrel{\rightharpoonup}{\mathrm{i}} \\ & \text { N } \end{aligned}$ | $\begin{aligned} & \text { ஸ } \\ & \underset{\sim}{\circ} \end{aligned}$ | $\begin{aligned} & \underset{\sim}{2} \\ & \underset{\sim}{2} \end{aligned}$ | $\underset{\sim}{i}$ | $\begin{aligned} & \infty \\ & \infty \\ & \stackrel{\infty}{-} \end{aligned}$ | $\begin{aligned} & \infty \\ & \infty \\ & 0 \end{aligned}$ |
|  | $\begin{aligned} & \text { S } \\ & \sum_{\mathbf{M}}^{\text {n }} \end{aligned}$ | $\begin{aligned} & \text { 品 } \\ & \hline \end{aligned}$ | $\begin{aligned} & \text { n } \\ & \underset{\sim}{+} \\ & + \\ & \stackrel{n}{\infty} \\ & \underset{\sim}{0} \end{aligned}$ | $\begin{gathered} \text { 上 } \\ +1 \\ + \\ \underset{N}{\mathrm{~N}} \end{gathered}$ | $\begin{aligned} & \stackrel{0}{e} \\ & + \\ & + \\ & \stackrel{0}{6} \\ & \hline \end{aligned}$ |  | $\begin{aligned} & 2 \\ & + \\ & + \\ & 0 \\ & 0 \\ & \mu \end{aligned}$ | $\begin{gathered} \infty \\ + \\ 0 \\ 0 \\ \end{gathered}$ | $\begin{aligned} & \mathrm{I} \\ & + \\ & + \\ & \sim \\ & \infty \\ & i n \\ & \hline \end{aligned}$ |
|  | － | 7 | $\stackrel{\square}{1}$ | 9 | 7 | $\overrightarrow{1}$ | $\bigcirc$ | ¢ | 7 |
|  | $0$ |  | $\underset{\text { oin }}{+}$ |  | $\begin{gathered} 1 \\ \text { B/w } \end{gathered}$ | $\infty$ |  | $+$ | $\square_{010}^{1}$ |
|  | H | $\stackrel{N}{\mathrm{~V}}$ | － | $\cdots$ | $\stackrel{\mathrm{V}}{ }$ | 7 | $\cdots$ | －10 | － |
|  |  | ＋ | $\bigcirc$ | $\bigcirc$ | $\bigcirc$ | ＋ |  | 10 | － |
| 首 |  |  | ${ }_{\sim}^{*}$ |  |  | $\stackrel{*}{*}$ |  |  | ${ }^{\circ}$ |
| งsยปว |  | səłełS эฺ̣uorədКH |  |  |  |  |  |  |  |


Mesonic resonant states, November 1963

| $\begin{array}{\|c} \text { Sym- } \\ \text { bol } \end{array}$ | Charge | Quantum numbers |  |  | Mass |  | $\left\lvert\, \begin{gathered} \text { Full } \\ \text { width } \Gamma \\ (\mathrm{MeV}) \end{gathered}\right.$ | $\left.\begin{gathered} \text { Life- } \\ \text { time } \\ \Gamma^{-1} \\ \left(1 m_{\pi \pm}\right) \end{gathered} \right\rvert\,$ | Production |  | Decay |  |  | Footnotes |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  | $\begin{array}{\|l\|l} \text { Pro- } \\ \text { cess } \end{array}$ | $\begin{aligned} & k_{1 \mathrm{ab}} \\ & (\mathrm{MeV}) \end{aligned}$ |  |  | Modes | Branching ratio (\%) | $\begin{gathered} Q \\ (\mathrm{MeV}) \end{gathered}$ |  |
|  |  | $T$ | $J \quad P \quad G S$ | $S$ |  |  |  |  |  |  |  | ( MeV ) | $\left(m_{\pi \pm}\right)$ |  |
|  | 0 | $\leqslant 1$ |  | 0 | 1410 | 10.1 | 50-60 | 2.5 | $\overline{\mathrm{p}} \mathrm{p}$ |  | $\mathrm{K}_{1}^{\mathrm{o}} \mathbf{K}^{ \pm} \pi^{\mp}$ |  | 280 | 37 |
| f | 0 | 0 | $2++$ | 0 | $1256 \pm 14$ | 9.0 | 160 | 0.88 | $\pi^{-p}$ | 2070 | $\pi^{-} \pi^{+}$ | 100 | 970 | 38 |
| $\begin{aligned} & \mathrm{B} \\ & \mathrm{~B} \\ & \hline \end{aligned}$ | + | 1 | + | 0 | $\begin{aligned} & 1215 \\ & 1220 \end{aligned}$ | $\begin{aligned} & 8.7 \\ & 8.7 \\ & \hline \end{aligned}$ | $\begin{array}{\|c} \hline 170 \\ 100 \pm 20 \\ \hline \end{array}$ | $\begin{aligned} & 0.8 \\ & 1.4 \end{aligned}$ | $\begin{aligned} & \pi^{-} \mathrm{p} \\ & \pi^{+} \mathrm{p} \end{aligned}$ | $\begin{aligned} & 1990 \\ & 1990 \end{aligned}$ | $\begin{aligned} & \pi-\omega \\ & \pi^{+} \omega \end{aligned}$ |  | $\begin{aligned} & 295 \\ & 300 \end{aligned}$ | $\begin{aligned} & 39 \\ & 40 \end{aligned}$ |
|  | 0 | 0 | even + + 0 |  | $1040 \pm 40$ | 7.4 |  |  | $\mathrm{K}^{-} \mathrm{p}$ | 1780 | $\begin{aligned} & \mathrm{K}_{1}^{0} \mathrm{~K}_{1}^{0} \\ & 2 n \times \pi \end{aligned}$ |  | 44 | 41 |
| $\phi$ | 0 | 0 | 1 - - |  | $1019.5 \pm 0.3$ | 7.30 | $3: 1 \pm 0.6$ | 45.0 | $\mathbf{K}^{-p}$ | 1760 | $\begin{aligned} & \mathrm{K}_{1}^{0} \mathrm{~K}_{2}^{\mathrm{o}} \\ & \left(\begin{array}{l} 2 n+1) \times \pi \\ \rho \pi \end{array}\right. \end{aligned}$ | $\begin{aligned} & 90 \pm 10 \\ & 10 \pm 10 \end{aligned}$ | $\begin{array}{r} 24 \\ 130 \end{array}$ | 42 |
|  | - | $\geqslant 1$ |  | 0 | $1000 \pm 10$ | 7.1 | $120 \pm 30$ | 1.2 | $\pi^{-p}$ | 1515 | $\pi^{-} \pi^{0} \pi^{0}$ |  | 590 | 43 |
|  | 0 | 0 |  | 0 | $922 \pm 30$ | 6.6 | <150 | $>0.9$ | $\pi^{-p}$ | 1360 | $\pi^{-} \pi^{+}$ |  | 643 | 44 |
|  | - | $\frac{1}{2}$ |  |  | $890.4 \pm 1.2$ | 6.38 | $51 \pm 2$ | 2.74 | $\begin{aligned} & K^{-} p \\ & K^{-} p \\ & \pi^{-}-p \\ & \pi^{-}-p \end{aligned}$ | $\begin{aligned} & 1074 \\ & 1078 \\ & 1834 \\ & 1657 \end{aligned}$ | $\begin{aligned} & \overline{\mathbf{K}}^{\circ} \pi^{-} \\ & \mathbf{K}^{-} \mathrm{o} \end{aligned}$ | $\begin{aligned} & 60 \pm 16 \\ & 40 \pm 16 \end{aligned}$ | 254 | 45 |
| $\overline{\mathrm{K}}$ * | 0 |  |  |  | $\begin{aligned} & \mathbf{K}^{-r} \pi^{+} \\ & \frac{\mathrm{K}^{0}}{\pi} 0 \end{aligned}$ |  |  |  |  |  |  | 258 |  |  |
| K* | + |  |  |  | $\mathrm{K}^{\mathrm{K}} \mathrm{m}^{\text {O}}$ |  |  |  |  |  | 67 | 259 |  |  |
|  |  |  |  |  | $\mathrm{K}^{+} \mathrm{m}^{0}$ |  |  |  |  |  | 33 | 263 |  |  |
| K* | 0 |  |  |  | $\left\{\begin{array}{l} \mathbf{K}^{+} \pi- \\ \mathbf{K}^{\circ} \pi \end{array}\right.$ |  |  |  |  |  |  | $\begin{aligned} & 258 \\ & 259 \end{aligned}$ |  |  |
|  |  |  |  |  | $\mathrm{K}^{\circ} \mathrm{O}$ |  |  |  |  |  |  |  |  |  |


| $\omega$ | 0 | 0 |  | $1-\ldots$ |  | 783 | $\pm 2$ | 5.6 | $9.5 \pm 2.1$ | 14.7 | $\overline{\mathrm{p}} \mathrm{p}$ |  | $\left\{\begin{array}{l} \pi^{+} \pi^{-} \pi^{o} \\ \pi o \gamma \\ \pi^{+} \pi^{-} \\ \mathrm{e}^{+} \mathrm{e}^{-} \\ \pi^{+} \pi^{-} \gamma \end{array}\right.$ | $\begin{array}{c\|} 85 \\ 9.6 \pm 1.2 \\ 3.9 \pm 1.4 \\ 0.4 \pm 0.26 \end{array}$ | $\begin{aligned} & 368 \\ & 647 \\ & 503 \\ & 782 \\ & 503 \end{aligned}$ | 46 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\rho$ | - | 1 |  | $1-+$ |  | $757$ | $\pm 5$ | 5.4 | $120 \pm 10$ | $1.2$ | $\pi^{-p}$ | $1029$ | $\left\|\begin{array}{l} \pi^{-} \pi^{0} \\ \pi^{-} \pi^{0} \pi^{\circ} \\ \pi^{-\pi} \mathrm{m}_{\pi} \mathrm{o} \mathrm{o} \\ \pi^{-} \pi^{+} \pi^{-} \pi^{\mathrm{o}} \\ \pi \gamma \end{array}\right\|$ | $>86$ $<3$ $\leqslant 4$ $<2$ $<5$ | $\begin{aligned} & 475 \\ & 340 \\ & 205 \\ & 196 \\ & 617 \end{aligned}$ | 47 |
| $\rho$ <br> $\rho$ | 0 + + |  |  |  | 0 | 754 | $\pm 5$ | 5.4 | $110 \pm 10$ | 1.3 | $\begin{aligned} & \pi \mathrm{N} \\ & \pi^{+} \mathrm{p} \end{aligned}$ | $\begin{gathered} 1029 \\ 1066 \end{gathered}$ | $\left\|\begin{array}{l} \pi^{-} \pi^{+} \\ \text {neutrals } \\ \pi^{+} \pi^{-} \pi^{+} \pi^{-} \\ \pi^{+} \pi^{\circ} \end{array}\right\|$ | $\begin{gathered} 94(+6 /-40) \\ 6(+40 /-6) \\ <2 \end{gathered}$ | $\begin{aligned} & 470 \\ & 191 \\ & 495 \end{aligned}$ | 48 |
| $\begin{aligned} & \kappa \\ & \kappa \end{aligned}$ | 0 | $\frac{1}{2}$ | $\geqslant$ | 1 |  | 725 | $\pm 2$ | 5.2 | <12 | >12 | $\pi^{-} \mathrm{p}$ | 1485 | $\begin{aligned} & \hline \mathbf{K}^{+} \pi^{-} \\ & \mathbf{K}^{0} \pi^{\mathrm{o}} \\ & (\mathbf{K} \pi)^{+} \\ & \hline \end{aligned}$ |  | $\begin{array}{\|c\|} \hline 96 \\ 97 \\ 92-101 \\ \hline \end{array}$ | 49 |
| $\eta$ | 0 | 0 |  | $0-+$ |  | 548.5 | $\pm 0.6$ | 3.93 | $\leqslant 7$ | $\geqslant 20$ | $\pi \sim p$ | 685 | $\begin{array}{\|l\|} \hline \pi^{+} \pi^{-} \pi^{0} \\ \pi^{+} \pi^{-} \gamma \\ 3 \pi^{0} \\ \pi^{0} \gamma \gamma \\ 21 \\ 2 \gamma \end{array}$ | $\left.\begin{array}{c} 25 \pm 8 \\ 7 \pm 2 \\ 1 \pm \pm 9 \\ 16 \pm 16\} 68 \pm 9 \\ 31 \pm 17 \end{array}\right\}$ | $\begin{aligned} & 135 \\ & 270 \\ & 144 \\ & 414 \\ & 549 \\ & \hline \end{aligned}$ | 50 |
|  | 0 | 0 |  |  | 0 | 520 | $\pm 20$ | 3.7 | $70 \pm 30$ | 2.0 | $\pi^{-p}$ | 639 | $\pi^{+} \pi^{-}$ |  | 240 | 51 |
| $\omega_{\mathrm{ABC}}$ | 0 | 0 |  | $0++$ |  | 317 | $\pm 6$ | 2.3 | $\leqslant 16$ | $\geqslant 9$ | pd |  | $\pi^{+} \pi^{-}$ |  | 38 | 52 |

## FOOTNOTES

1. $m\left(\Xi^{-}\right)$is a weighted average of $1321.1 \pm 0.65$ Schneider (1963)
$1321.4 \pm 0.4$ Jauneau et al. (1963a)
$1311.1 \pm 0.2$ Connolly et al. (1963c).
$\tau\left(\Xi^{-}\right)$is a weighted average of the following measurements in $10^{-10} \mathrm{~s}$ :
1.86 (+0.15/-0.14) Jauneau et al. (1963a)
$1.55 \pm 0.31 \quad$ Schneider (1963)
1.74 ( $+0.18 /-0.15$ ) Connolly et al. (1963c)
$1.75 \pm 0.07$ L. Alvarez, J. Berge, J. Hubbard, R. Kalbfleisch, J. Sha-
fer, F. Solmitz, M. Stevenson and S. Wojcicki, unpublished work quoted by Ticho (1963)
$1.77 \pm 0.12$ D. Carmony, G. Pjerrou, P. Schlein, W. Slater, D. Stork, H. Ticho, unpublished work quoted by Ticho (1963).
$J=\frac{1}{2}$ is compatible with the data of Bertanza et al.
2. $\quad m\left(\Xi^{0}\right)$ is obtained from $m\left(\Xi^{-}\right)$and the weighted average $\sigma .0 \pm 1.0$ of the following measurements of the mass difference $m\left(\Xi^{-}\right)-m\left(\Xi^{0}\right)$
$3.6 \pm 2.4$ Connolly et al. (1963c)
$6.8 \pm 1.5$ Jauneau et al. (1963b)
6.1 $\pm$ 1.6 D. Carmony, G. Pjerrou, P. Schlein, W. Slater, D. Stork, H. Ticho, unpublished work quoted by Ticho (1963).
$\tau\left(\Xi^{0}\right)$ is the weighted average given by Ticho (1963).
3. $m\left(\bar{\Xi}^{0}\right)$ is reported by Baltay et al. (1963b).
4. The $\Sigma$ masses are from Barkas et al. (1963).
$\tau\left(\Sigma^{+}\right)$and $\tau\left(\Sigma^{-}\right)$are weighted averages of the following results in $10^{-10} \mathrm{~s}$ :
$\left.\begin{array}{l}\tau\left(\Sigma^{+}\right)=0.81(+0.06 /-0.05) \\ \tau\left(\Sigma^{-}\right)=1.61(+0.10 /-0.09)\end{array}\right\}$ Barkas et al. (1960)
$\left.\tau\left(\Sigma^{-}\right)=1.61(+0.10 /-0.09)\right\}$ Barkas et al. (1960)
$\left.\begin{array}{l}\tau\left(\Sigma^{+}\right)=0.765 \pm 0.04 \\ \tau\left(\Sigma^{-}\right)=1.58 \pm 0.06\end{array}\right\}$ Humphrey et al. (1962).
The $\Sigma^{+}$branching ratio is a weighted average of
$51.0 \pm 2.4$ Humphrey et al. (1962)
$48 \pm 7$ Granet (1962).
The $\Sigma$ parity is + from Tripp et al. (1962), who report $P(\Sigma \mathrm{pK})=-1$.
5. The upper limit of $\tau\left(\Sigma^{0}\right)$ is from Alvarez et al. (1957); the lower limit from Dreitlein et al. (1961), who also give a theoretical estimate

$$
\tau\left(\Sigma^{0}\right)=1.1 \times 10^{-19} \mathrm{~s}
$$

6. $m(\Lambda)$ is a weighted average of
$1115.36 \pm 0.14$ Barkas et al. (1960)
$1115.46 \pm 0.15$ Bhowmik et al. (1961)
$1115.25 \pm 0.36$ Armenteros et al. (1962b)
$1115.04 \pm 0.41$ Baltay et al. (1962).
$\mu(\Lambda)$ is the value of Cool et al. (1962), which disagrees with the value $0.0 \pm 0.6$ nm , reported by Kernan et al. (1963).
$\tau(\Lambda)$ is a central value from the ideogram of Crawford (1962). The new values which have appeared since 1962 do not resolve the so called East-West effect.
$m(\bar{\Lambda})$ is a weighted average of
$1115.40 \pm 0.39$ Baltay et al. (1962)
$1115.52 \pm 0.55$ Armenteros et al. (1962b).
$\tau(\bar{\Lambda})$ is from Baltay et al. (1962).
The $\Lambda \rightarrow \mathrm{p} \pi^{-}$branching ratio is a central value of
$64.3 \pm 1.6$ Humphrey et al. (1962)
$64.5 \pm 2.2$ Crawford (1962), weighted average based on earlier measurements; $68.5 \pm$ 1.7 J. Anderson, F. and B. Crawford, R. Golden, L. Lloyd, G. Meisner and L. Price (unpublished), quoted by Crawford (1962)
70.9 $\pm$ 3.4 Chrétien et al. (1963).

The $\Lambda$ parity is + from Cronin et al. (1962), who give support for $P(\mathrm{~K} \Lambda)=-1$.
7. The nucleon and electron masses are from R. Cohen et al. (1957).
$\mu(\mathrm{n})$ is from V. Cohen et al. (1956).
The neutron mean life is based on the value $11.7 \pm 0.3 \mathrm{~min}$ for the half life, by Sosnovskij et al. (1959).
8. $\mu(\mathrm{p})$ is a weighted average of the following data in nuclear magnetons:
$2.79277 \pm 0.00005$ Sanders et al. (1963a)
$2.792765 \pm 0.000060$ Sommer et al. (1951)
$2.792810 \pm 0.000076$ average value computed using the electron and proton masses and the following values of $\mu(\mathrm{p})$ in Bohr magnetons:
$657.442 \pm 0.003$ Liebes et al. (1959)
$657.4436 \pm 0.0025$ Sanders et al. (1963b).
9. Button et al. (1962b).
10. $m\left(\mathrm{~K}^{+}\right)$is from Barkas et al. (1963).
$\tau\left(\mathrm{K}^{+}\right)$is a weighted average of the following results in $10^{-8} \mathrm{~s}$ :
$1.224 \pm 0.013$ Barkas et al. (1960)
$1.231 \pm 0.011$ Boyarski et al. (1962).
The $\mathrm{K}^{+}$branching ratios are from Roe et al. (1961) and Giacomelli et al. (1963). Note that these branching ratios (except for $\tau$ ) disagree with the weighted averages obtained from emulsion experiments, as quoted by Crawford (1962):

$$
\begin{aligned}
\mu 2 & =57.4 \pm 2.0 \\
\pi 2 & =25.6 \pm 1.5 \\
\mu 3+\mathrm{e} 3+\tau & =11.0 \pm 1.0 \\
\tau & =5.7 \pm 0.2
\end{aligned}
$$

11. $\quad m\left(\mathrm{~K}^{0}\right)$ is from Rosenfeld et al. (1959). $\mu\left(\mathrm{K}^{0}\right)$ is from Okonov (1962). $\tau\left(\mathbf{K}_{1}^{\mathbf{0}}\right)$ is a weighted average, based on the following recent results only (in $10^{-10} \mathrm{~s}$ ), all quoted by Crawford (1962):
$0.94 \pm 0.05$ Crawford et al. (1962)
$0.90 \pm 0.05$ A. F. Garfinkel, Report Nevis 104 (1962) (Thesis, Columbia University Physics Department),
$0.885 \pm 0.025$ R. L. Golden, G. Alexander, J. A. Anderson, F. S. and B. B. Crawford, L. J. Lloyd, G. W. Meisner, and L. Price (to be published).

The $K_{1}^{0} \rightarrow \pi^{0} \pi^{0}$ branching ratio is from Brown et al. (1963). The values $28.5 \pm 3.6$ by Crawford et al. (1959), $28.8 \pm 2.1$ by Chrétien et al. (1963), $26.0 \pm 2.4$ by Anderson et al. (cf. reference 6), and $29 \pm 3$, weighted average of earlier results, reported at the 1960 Rochester Conference, are less accurate and do not agree with theory.
12. $\tau\left(\mathrm{K}_{2}^{\mathbf{0}}\right)$ is a weighted average of the following results in $10^{-8} \mathrm{~s}$ :
6.8( $+2.6 /-1.5$ ) Crawford (1962)
8.1 (+3.3/-2.4) Bardon et al. (1958)
5.1(+2.4/-1.3) Darmon et al. (1962).

The branching ratios have been obtained from the following results:
$R_{1}=\frac{\mathbf{K}_{2}^{0} \rightarrow 3 \pi^{0}}{\mathbf{K}_{2}^{0} \rightarrow \text { all charged }}=0.38 \pm 0.07$ Anikina et al. (1962)
$R_{\mathbf{2}}=\frac{\mathbf{K}_{2}^{0} \rightarrow \pi^{+} \pi^{-} \pi^{0}}{\mathbf{K}_{\mathbf{2}}^{0} \rightarrow \text { all charged }}=0.127 \pm 0.020$ Luers et al. (1961)
$R_{2}=0.134 \pm 0.018$, Anikina et al. (1962)
$R_{3}=\frac{\mathbf{K}_{2}^{0} \rightarrow \pi \mathrm{e} v_{\mathrm{e}}}{\mathrm{K}_{2}^{0} \rightarrow \text { all charged }}=0.458 \pm 0.048$, Luers et al. (1961)
$R_{3}=0.415 \pm 0.120$, Astier et al. (1961).
A value on $R_{3}=0.185(+0.038 /-0.034)$ by Astier et al. has not been used.
13. $m\left(\pi^{ \pm}\right)$is from Shapiro et al. (1962).
$\tau\left(\pi^{ \pm}\right)$is from Merrison (1962).
14. $\tau\left(\pi^{0}\right)$ is from von Dardel et al. (1963).

The $\pi^{0}$ mass has been obtained from the $\pi^{+}$mass and the measurement by Czirr (1963) of the mass difference $m\left(\pi^{-}\right)-m\left(\pi^{0}\right)=(4.6056 \pm 0.0055) \mathrm{MeV}$.
The relative frequency for $\gamma \mathrm{e}^{+} \mathrm{e}^{-}$-decay is from J. Tietge et al. (1962).
15. $\mu\left(\mu^{+}\right)$is from Charpak et al. (1962).
$m\left(\mu^{+}\right)$is a combined value of the latest measurements as given by G. Bingham (1963).
$\tau\left(\mu^{+}\right)$is a central value including all measurements quoted by Lundy (1962) and the values
$\left.\begin{array}{l}\tau\left(\mu^{+}\right)=(2.197 \pm 0.002) \mu \mathrm{s} \\ \tau\left(\mu^{-}\right)=(2.198 \pm 0.002) \mu \mathrm{s}\end{array}\right\} \quad$ Meyer et al. $(1963)$
$\tau\left(\mu^{+}\right)=(2.202 \pm 0.004) \mu \mathrm{s} \quad$ Eckhause et al. (1963).
$\mu(\mathrm{e})$ is from Schupp et al. $(1961)$.
17. Bahcall et al. (1961).
18. Langer et al. (1952).
19. $v_{\mu}$ and $v_{\mathrm{e}}$ are left-handed screw states, $\bar{v}_{\mu}$ and $\bar{v}_{\mathrm{e}}$ right-handed, $\gamma$ has two spin states: right- and left-handed. We define a left-(right)-handed screw state as a state with negative (positive) spin.
20. This possible resonance was found in the $\Sigma^{ \pm} \pi^{+} \pi^{-}$channels but not in the $\Sigma^{ \pm} \pi^{ \pm} \pi^{ \pm}$or $\Sigma^{\mp} \pi^{ \pm} \pi^{ \pm}$channels, using $3.5 \mathrm{GeV} / c \mathrm{~K}^{-}$-mesons, by the Glasgow, Imperial College, Oxford and Rutherford Laboratory (1963).
21. Quantum numbers, mass and width values are taken from Barbaro-Galtieri et al. (1963b). Further evidence has been given by Kerth (1961), Chamberlain et al. (1962), Kuznetsov et al. (1962a) and Beall et al. (1962, 1963).
22. Belliere et al. (1963), Halsteinslid et al. (1963).
23. Barbaro-Galtieri et al. (1963b). Some previous indications have been given by Belyakov et al. (1962b) and Grashin et al. (1963).
24. The $\overline{\mathbf{K}} \mathbf{N}$ branching ratio is from Smith (1963b): An indication for positive parity has been given by Taher-Zadeh et al. (1963), all other data are from Alvarez et al. (1963). For further evidence, cf. reference list in Roos (1963).
25. Indications by March et al. (1962), Koch et al. (1962) and Baltay et al. (1963a). Evidence against has been given by Kalbfleisch et al. (1963).
26. $m\left(\Xi_{13}^{*}\right)$ is a weighted average of the following results in MeV :
$1529 \pm 5$ Pjerrou et al. (1962)
$1535 \pm 3$ Bertanza et al. (1962b).
The spin and parity assignments, the width and the $\Xi_{13}^{* 0}$ branching ratios are from Schlein et al. (1963). Note that this width is in disagreement with the value $\Gamma=16 \pm 3 \mathrm{MeV}$, given by Connolly et al. (1963b).
27. $m\left(\mathrm{Y}_{03}^{*}\right)$ is a weighted average of
$1519.4 \pm 2$ Watson et al. (1963)
$1517.2 \pm 3$ Barbaro-Galtieri et al. (1963b).
The width and branching ratios are from Watson et al. (1963).
28. $\quad m\left(\mathrm{Y}_{0}^{*}\right)$ is from Frisk et al. (1962).

This mass value is confirmed by Samman et al. (1963) who give ( $1405.1 \pm 0.9$ ) MeV . The width claimed by Frisk et al. (1962) has, however, not been confirmed in similar experiments by Barbaro-Galtieri et al. (1963a) nor by Lokanathan
et al. (1963). The width of Samman et al. (1963) is 10.3 MeV , or intermediate to that of Frisk et al. (1962) and the commonly accepted bubble chamber value, which we have listed.

Determinations of the spin and the parity are consistent with $J^{P}=\frac{1}{2}^{-}$, but this assignment is not finally settled, cf. Dalitz (1963).
29. $m\left(\mathrm{Y}_{13}^{*}\right)$ is a weighted average of
$1382 \pm 3$ Dahl et al. (1961)
$1385 \pm 5$ Alston et al. (1961a)
$1380 \pm 3 \quad$ Bertanza et al. (1963)
$1383.5 \pm 4$ Cooper et al. (1963), $\mathrm{Y}_{13}^{*+}$ mass
$1381 \pm 4 \quad$ Curtis et al. (1963).
The value $1389 \pm 3$ of Baltay et al. (1963c) has not been used.
The spin and parity assignments are from Shafer et al. (1963) and Bertanza et al. (1963). $\Gamma\left(\mathrm{Y}_{13}^{*}\right)$ is a weighted average of
$50 \pm 10$ Bertanza et al. (1963)
$30 \pm 9$ Curtis et al. (1963).
The value $26 \pm 5$ of Baltay et al. (1963c) has not been used. The branching ratios are from Alston et al. (1961a).
30. Lander et al. (1963).

Some indications at 2.6 GeV have been reported by the Aachen-Berlin-Bir-mingham-Bonn-Hamburg-London (I.C.)-München collaboration (1963).
31. Diddens et al. (1962). An indication for $J\left(\mathrm{~N}_{1}^{*}\right)=\frac{9}{2}$ is reported by Alexander et al. (1963).
32. Klepikov et al. (1960), Falk-Vairant et al. (1961), Moyer (1961), Devlin et al. (1962), Helland et al. (1963), Salin (1963). The K $\Sigma$-decay mode has been reported by March et al. (1962) and W. Walker (1962b), and the K $\Lambda$-decay mode by Erwin et al. (1962a), Kuznetsov et al. (1962b) and W. Walker et al. (1963b). The branching ratios are from Rosenfeld (1963).
33. The mass value is obtained from the weighted average of
$T_{\pi}=(890 \pm 9) \mathrm{MeV}$ Falk-Vairant et al. (1961)
$T_{\pi}=(900 \pm 15) \mathrm{MeV}$ Devlin et al. (1962).
The width is given by Omnès et al. (1961). The branching ratios are from A. H. Rosenfeld (1963). Further possible evidence for $\mathrm{K} \Lambda$ decay has been reported by Kuznetsov et al. (1962b), Baz et al. (1962) and Bertanza et al. (1962a).

For further details, cf. Klepikov et al. (1960), Moyer et al. (1961), Feld et al. (1962), Helland et al. (1963), Salin (1963), R. Walker (1963) and Cocconi et al. (1963).
34. The mass value is obtained from the weighted average of $T_{\pi}=(605 \pm 5) \mathrm{MeV}$ Falk-Vairant et al. (1961)
$T_{\pi}=(600 \pm 15) \mathrm{MeV}$ Devlin et al. (1962)
$E_{\gamma}=(740 \pm 10) \mathrm{MeV}$ Deutsch et al. (1961)
$E_{\gamma}=(750 \pm 15) \mathrm{MeV}$ Bellettini et al. (1963),
where $T_{n}=E_{\gamma}-140 \mathrm{MeV}$.
$\Gamma\left(\mathrm{N}_{13}^{*}\right)$ is a central value, which includes the values of Omnès et al. (1961), Devlin et al. (1962), Feld et al. (1962), Salin (1963) and Cocconi et al. (1963), but does not agree with the value $<60 \mathrm{MeV}$ given by Deutsch et al. (1961). See also Klepikov et al. (1960).
35. The $\gamma p$ evidence has been communicated privately from the Stoppini-group at Frascati by R. Gomez. The pp indication has been reported by Cocconi et al. (1963). Some further indications of a resonance at 1480 MeV , which may be a mixture of $\mathrm{N}_{13}^{*}$ and the 1400 MeV resonance, has been reported by Pauli et al. (1963).
36. The mass value is from Klepikov et al. (1960) and corresponds to the position where the phase shift passes through $90^{\circ}$ with increasing energy, cf. also Dalitz (1963). The position of the maximum seems to be located at about ( $1220 \pm 10) \mathrm{MeV}$, cf. Klepikov et al. (1960), Hart et al. (1962), Samios et al. (1962), Duboc et al. (1963), Chadwick et al. (1963b) and R. Walker (1963).

The width is a weighted average of

$$
90 \pm 20 \text { Samios et al. (1962) }
$$

$100 \pm 25$ Chadwick et al. (1963b), in good agreement with Klepikov et al. (1960), Detoeuf (1961) and R. Walker (1963).

For further details, cf. Falk-Vairant et al. (1961), Moyer (1961), Feld et al. (1962) and Salin (1963).
37. Armenteros et al. (1963).
38. $m\left(\mathrm{f}^{0}\right)$ is a weighted average of
$1260 \pm 35$ Veillet et al. (1963)
$1250 \pm 25$ Selove et al. (1962)
$1250 \pm 50$ Guiragossian (1963)
$1260 \pm 20$ Bondar et al. (1963a).
$\Gamma\left(\mathrm{f}^{0}\right)$ is from Bondar et al. (1963a).
39. Chung et al. (1963). An indication has been reported by Bondar et al. (1963b).
40. Abolins et al. (1963).
41. Alexander et al. (1962a), Erwin et al. (1962b), H. Bingham et al. (1962), Bigi et al. (1962), A. H. Rosenfeld (1963), Armenteros et al. (1963). Possible $4 \pi$ decays have been reported by Xuong et al. (1962a, b), Chadwick et al. (1962), Foelsche et al. (1962a) and Lynch (1962).
42. $m(\phi)$ is a weighted average of
$1018.6 \pm 0.5$ Gelfand et al. (1963b)
$1020.5 \pm 0.5$ Connolly et al. (1963a)
$1019 \pm$ 1.6 Glasgow, Imperial College, Oxford and Rutherford Laboratory (1963).
$\Gamma(\phi)$ is a weighted average of
$3.1 \pm 1.0$ Gelfand et al. (1963b)
$3.1 \pm 0.8$ Connolly et al. (1963a).
The $\rho \pi$ branching ratio is from Connolly et al. (1963a) who also put an upper limit of $8 \%$ to the $\pi \pi / K \bar{K}$ branching ratio.
43. Trebukhovsky et al. (1963). Some earlier indications have been reported by Ainutdinov et al. (1962) in the $T=2$ channels ( $\pi^{ \pm} \pi^{ \pm}$) and by Guiragossian et al. (1962) in the ( $\pi^{+} \pi^{-}$)-channel.
44. Hulubei et al. (1963). Erwin et al. (1963) give $T=0, m=940 \mathrm{MeV}$ and $\Gamma \sim 20 \mathrm{MeV}$.
45. $m\left(\mathrm{~K}^{*}\right)$ and $\Gamma\left(\mathrm{K}^{*}\right)$ are weighted averages of the following results in MeV :
$m\left(\mathrm{~K}^{*}\right) \quad \Gamma\left(\mathrm{K}^{*}\right)$
890.4( $\pm 2) \quad 47( \pm 5) \quad$ Alston et al. (1962)

885 ( $\pm 5$ ) $60 \pm 5$ Alexander et al. (1962b)
$897 \quad \pm 10 \quad 60 \pm 10$ Colley et al. (1962)
$885 \pm 5 \quad 55 \pm 5$ Armenteros et al. (1962c)
885 ( $\pm 10$ ) 50( $\pm 10)$ Smith et al. (1963a)
$898 \quad \pm 5 \quad 46 \pm 8$ Chadwick et al. (1963b)
$892 \pm 2 \quad 50 \pm 5$ Kraemer et al. (1963)
$888 \pm 3 \quad 45 \pm 5$ Gelsema et al. (1963)
$35 \pm 15$ March et al. (1962).
The spin and parity assignments are from Chinowsky et al. (1962).
The $\mathrm{K}_{1}^{*-}$ branching ratios are weighted averages of the following results:
$R=\Gamma\left(\overline{\mathrm{K}}^{0} \pi^{-}\right) / \Gamma\left(\mathrm{K}^{-} \pi^{0}\right)=1.4 \pm 0.4$ Alston et al. (1961b)
$1 / R=0.5 \pm 0.2$ Graziano et al. (1962).
The $\pi^{-}$momentum for $K_{1}^{* 0}$ production is here tabulated as $1657 \mathrm{MeV} / c$ which corresponds to associated production with $\Lambda$. Associated production with $\Sigma^{0}$ gives $1826 \mathrm{MeV} / c$.
46. $m\left(\omega^{0}\right)$ is a central value of
$779.4 \pm 1.4$ Armenteros et al. (1962a)
$784.0 \pm 0.9$ Gelfand et al. (1963a).
$\Gamma(\omega)$ is from Gelfand et al. (1963a).
The ( $\omega \rightarrow$ neutrals) $/(\omega \rightarrow 3 \pi$ ) branching ratio is a weighted average of
$0.135 \pm 0.035$ Steinberger et al. (1963)
$0.11 \pm 0.03$ Murray et al. (1963)
$0.11 \pm 0.02$ Buschbeck-Czapp et al. (1963)
$0.09 \pm 0.04$ Fields et al. (1963).
The branching ratio ( $\omega \rightarrow 2 \pi$ )/( $\omega \rightarrow 3 \pi$ ) is from Murray et al. (1963), in agreement with other measurements: W. Walker (1963a), Fickinger (1963) and the value $5 \%$ at 2.75 GeV communicated privately by Puppi, and in only mild disagreement with Steinberger et al. (1963).

The branching ratio $\left(\omega \rightarrow \mathrm{e}^{+} \mathrm{e}^{-}\right) /(\omega \rightarrow$ neutrals $)=(0.5 \pm 0.3) \%$ has been quoted by Berthelot (1963). According to Barmin et al. (1963) all the neutral decays consist of $\pi^{0} \gamma$.

The $\pi^{+} \pi^{-} \gamma$ decay has been seen by Belyakov et al. (1962a), cf. also Nguyen Dinh Tu (1962).
47. $m\left(\rho^{-}\right)$and $\Gamma\left(\rho^{-}\right)$are weighted averages of the following results, tabulated as functions of the lab kinetic energy $T_{\pi}$ :

| $\begin{aligned} & m(\rho \pm) \\ & (\mathrm{MeV}) \end{aligned}$ | $\begin{aligned} & \Gamma(\rho \pm) \\ & (\mathrm{MeV}) \end{aligned}$ | $\begin{gathered} T_{\pi} \\ (\mathrm{GeV}) \end{gathered}$ | References |
| :---: | :---: | :---: | :---: |
| $713 \pm 81$ | $31 \pm 143$ | 0.91 | Foelsche et al. (1962a) |
| $748 \pm 16$ |  | 1.1 | Kenney et al. (1962a, b) |
| $755 \pm 10$ | $61 \pm 24$ | 1.26 | Foelsche et al. (1962a) |
| $752 \pm 13$ |  | 1.45 | Saclay-Orsay-Bari-Bologna-collaboration (1961, 1963) |
| $740( \pm 13)$ | $120( \pm 15)$ | 1.72-1.93 | W. Walker et al. (1962a) |
| $755 \pm 10$ |  | 1.89-2.1 | W. Walker et al. (1963b) |
| $770 \pm 10$ | $130 \pm 10$ | 2.19-2.73 | Alff et al. (1962) |
| $775( \pm 25)$ | 125 $\pm$ (25) | 3.3 | Guiragossian (1963) |
| $780( \pm 25)$ |  | 7.0 | Grashin et al. (1962) |
| $755( \pm 15)$ | $110( \pm 15)$ | $\overline{\mathrm{p}} \mathrm{p}$ at rest | Chadwick et al. (1963a). |

The above mass values are all consistent with $757 \pm 5$, and no evidence for an energy dependence can be seen. In fact, the only such evidence has been presented by Foelsche et al., who give $m\left(\rho^{+}\right)=726 \pm 10$ and $\Gamma\left(\rho^{+}\right)=57 \pm 27$ at $T_{\pi}=1.09$ GeV . This value has not been used by us.

The upper limits of the branching ratios into $\pi^{-} \pi^{0} \pi^{0}$ and $\pi^{-} \pi^{+} \pi^{-} \pi^{0}$ have been given by Lynch (1962), into $\pi^{-} \pi^{0} \pi^{0} \pi^{0}$ by Alitti et al. (1962) and into $\pi \gamma$ by Berthelot (1963).
48. $m\left(\rho^{0}\right)$ and $\Gamma\left(\rho^{0}\right)$ are weighted averages of the following results, tabulated as functions of the lab kinetic energy $T_{\pi}$ :

| $m\left(\rho^{0}\right)$ <br> $(\mathrm{MeV})$ | $\Gamma\left(\rho^{0}\right)$ <br> $(\mathrm{MeV})$ | $T_{\pi}$ <br> $(\mathrm{GeV})$ | References |
| :---: | :---: | :--- | :--- |
| $752 \pm 27$ |  | 1.45 | Saclay-Orsay-Bari-Bologna collaboration (1963) |
| $760( \pm 15)$ |  | 1.72 | W. Walker et al. (1962a) |
| $750 \pm 10$ | $100 \pm 10$ | $2.19-2.73$ | Alf et al. $(1962)$ |
| $742( \pm 15)$ |  | 2.64 | Grashin et al. (1962) |
| $775( \pm 25)$ | $175( \pm 25)$ | 3.3 | Guiragossian $(1963)$ |
| $760 \pm 10$ | $90 \pm 10$ | $3.43-3.54$ | Abolins et al. (1963) |
| $747 \pm 17$ | $156 \pm 17$ | 4.55 | Samios et al. $(1962)$ |
| $750( \pm 15)$ |  | 7.0 | Grote et al. $(1962)$ |
| $755( \pm 15)$ | $110( \pm 15)$ | $\overline{\text { pp at rest }}$ | Kenney et al. $(1962 \mathrm{a}, \mathrm{b})$ |

There is no evidence for an energy dependence.
The branching ratios are from Meer et al. (1962).
The upper limit on the $\pi^{+} \pi^{-} \pi^{+} \pi^{-}$branching ratio has been given by Lynch (1962).
49. $\quad m(\kappa)$ is a weighted average of
$726 \pm 3$ Miller et al. (1963)
$723 \pm 3$ Wojcicki et al. (1963).
$\Gamma(\kappa)$ is from Wojcicki et al. (1963), in agreement with Miller et al. (1963), who give $\Gamma \leqq 20 \mathrm{MeV}$. Further support by Connolly et al. (1963b).
50. $m\left(\eta^{0}\right)$ is a weighted average of
$548 \pm 1$ Alff et al. (1962)
$548 \pm 1$ Foelsche et al. (1962a) at $1090 \mathrm{MeV} / \mathrm{c}$
$551 \pm 2$ Foelsche et al., at $1260 \mathrm{MeV} / \mathrm{c}$
$546 \pm 4$ Pickup et al. (1962)
$\left.\begin{array}{l}550 \pm 1.5 \\ 548 \pm 2\end{array}\right\}$ Bastien et al. (1962) from different decay modes.
$\Gamma\left(\eta^{0}\right)$ is from Bastien et al., Alff et al. and Foelsche et al. give $\Gamma<10 \mathrm{MeV}$.
The isospin assignment is from Carmony et al. (1962) and the spin and parities from Chrétien et al. (1962).

The branching ratios have been computed from the following measurements:
$R_{1}=\pi^{+} \pi^{-} \gamma / \pi^{+} \pi^{-} \pi^{0}=0.26 \pm 0.08$ Fowler et al. (1963)
$R_{2}=$ neutrals $/ \pi^{+} \pi^{-} \pi^{0}=2.5 \pm 0.5$, Alff et al. (1962)
$R_{2}=2.5 \pm 1.0$, Pickup et al. (1962)
$1 / R_{2}=0.31 \pm 0.11$, Bastien et al. (1962)
$R_{2}=3.0 \pm 0.7$, Fields et al. (1963)
$R_{2} /\left(1+R_{1}\right)=2.7 \pm 0.8$, Button-Shafer et al. (1962c)
$R_{2} /\left(1+R_{1}\right)=2.6 \pm 0.9$ Buschbeck-Czapp et al. (1963)
$R_{3}=2 \gamma /\left(3 \pi^{0}+\pi^{0} \gamma \gamma\right)=0.8 \pm 0.25$, Bacci et al. (1963)
$R_{3}=1.1 \pm 0.5$, Muller et al. (1963),
$\gamma \gamma /\left(\pi^{+} \pi^{-} \pi^{0}+\pi^{+} \pi^{-} \gamma\right)=0.99 \pm 0.48$, Crawford et al. (1963)
$3 \pi^{0} /\left(\pi^{+} \pi^{-} \pi^{0}+\pi^{+} \pi^{-} \gamma\right)=0.66 \pm 0.25$, Crawford et al. (1963).
51. Samios et al. (1962), Vittitoe et al. (1963), Kenney et al. (1963).
52. $m\left(\omega_{\mathrm{ABC}}\right)$ is a weighted average of
$310 \pm 10$ A. Abashian et al. (1960, 1961a, b)
$322 \pm 8$ B. Richter (1962).
Further evidence has been given by Button et al. (1962a) and Homer et al. (1963). The spin and parity assignments are from Abashian et al. (1963).

## References

The following shortenings have been used below:
Aix $=$ Proceedings of the Aix-en-Provence International Conference on Elementary Particles in 1961 (C.E.N. Saclay, Seine et Oise, France, 1962);
Athens $=$ Proceedings of the Athens Topical Conference on Recently Discovered Resonant Particles (Ohio University, Athens, Ohio, 1963);
CERN = Proceedings of the 1962 International Conference on High-Energy Physics at CERN (CERN, Geneva, 1962);
Sienna $=$ Proceedings of the Sienna International Conference on Elementary Particles (Frascati National Laboratories, Frascati, Italy, 1963)
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[^0]:    $\dagger$ "Not known" here and henceforth is short for "not known to"the compiler".

