# Tables of Elementary Particles and Resonant States 

Matts Roos<br>Nordisk Institut for T'eoretisk Atomfysik, Copenhagen, Denmark

## INTRODUCTION

THOSE reported particles which are customarily called elementary are contained in Table I.
Table II is a catalog, possibly incomplete, of those reported particles which are customarily called resonances. Many of these resonant states have not been generally accepted. We leave it, however, to the user of the table to put question marks against or delete such resonances which do not agree with his theory or experiment.

## CLASS, PARTICLE AND ANTIPARTICLE SYMBOLS

Division into hyperons, nucleons, and mesons is employed also for the resonances. The hyperonic states can then be logically defined as states with baryon number $B=1$ and strangeness $S \neq 0$, the nucleonic states as states with $B=1$ and $S=0$, and the mesonic states as states with $B=0$.

The leptons have been chosen so that $\mu^{-}, e^{-}, \nu_{\mu}$, and $\nu_{e}$ are particles; $\mu^{+}, e^{+}, \bar{\nu}_{\mu}$, and $\bar{\nu}_{e}$ antiparticles.

There are not yet name conventions for all resonances. For some of them many names are in use, out of which we have tried to choose the most convenient one. It is, for instance, not very practical to name the $\pi N$ resonances $N_{1}, N_{2}, N_{3}$, and $N_{4}$ or $N^{*}, N^{* *}, N^{* * *}$, and $N^{* * * *}$ now that a new $\pi N$ resonance has been found which comes in between $N_{3}$ and $N_{4}$ (or $N^{* * *}$ and $N^{* * * *}$ ).

We expect the starred notation to become unpopular by the time a resonance is discovered which needs eight stars to fit into the system.

Some resonances have not yet been named. Without any intention of depriving the discoverers of their right in this respect, we herewith use some working names. The working names have been formed using the following rules:

1. A few letters are used for particular classes of resonances:

$$
\begin{aligned}
& \Xi^{*} \text { for } S=-2, B=1 \\
& Y^{*} \text { for } S=-1, B=1 \\
& N^{*} \text { for } S=0, B=1 \text { when they decay into } \pi^{\prime} \text { 's } \\
& \text { and } N, \\
& Z^{*} \text { for } S=0, B=1 \text { when they decay into } \\
& \text { strange particles, } \\
& K^{*} \text { for } S= \pm 1, B=0
\end{aligned}
$$

$\kappa$ for bosons which decay into $K \bar{K}$-pairs, $\chi$ for bosons which decay into 4 pions,
$\psi$ for $T=2$ bosons which decay into 2 pions,
$\varphi$ for $T=0$ bosons which decay into 2 pions.
2. The baryons are given the two subscripts $2 T$, $2 J$ (for $\Xi^{*}, N^{*}$ and $Z^{*}$ ) or $T, 2 J$ (for $Y^{*}$ ), and the strange bosons $\left(K^{*}\right)$ one subscript, $2 T$. Not known subscripts are left blank. Further degeneracy in the notation is resolved by adding more stars.
3. $\kappa, \chi, \psi$ and $\varphi$ carry a subscript without physical significance, only to distinguish between different states.

Exceptions from rule 1 are $\omega_{\text {ABC }}$ (should be $\varphi_{0}$ ) and $f^{0}$ (should be $\varphi_{4}$ or $\psi_{6}$ ).

Clearly, symbols may change as further data become available or as other conventions are accepted.
In Table I particles and their found antiparticles are placed in separate columns. The only antiparticle not found (to our knowledge) is $\bar{\Xi}^{0}$. In Table II both particles and antiparticles are listed in the particle column.

## QUANTUM NUMBERS

A blank space in any of the quantum-number columns may signify that the quantity in question is not known, * or that it cannot be defined ( $T, T_{3}, S$, and parity for leptons). $T$ is not repeated for isospin multiplets, nor is $T, S$, or parity repeated for antiparticles. In Table II, $J$, parity and $G$ parity are not repeated for different charge states.

Different charge states are given separate entries when they have been found.

Parity is defined in relation to $N$ and $K$; by definition, $N$ has parity + and $K$ has - .

## MASS AND MAGNETIC MOMENT

The mass is given in two units, MeV and $m_{\pi^{ \pm}}$; for leptons $m_{\pi^{ \pm}}$is exchanged for $m_{\bullet}$. We choose the units

$$
\begin{aligned}
\hbar & =c=1 \\
m_{\pi^{ \pm}} & =139.58 \mathrm{MeV}=2.4881 \times 10^{-25} \mathrm{~g} \\
m_{e} & =0.510976 \mathrm{MeV}
\end{aligned}
$$

The accurate mass is expressed in MeV for all

[^0]Table I. Elementary Particles, March 1963.

|  |  |  | Isospin |  |  |  | Mass |  | Magnetic moment (e/2mp) | Mean life |  | Common decay modes | $\begin{array}{\|c\|} \hline \text { Branching } \\ \text { ratios } \\ (\%) \\ \hline \end{array}$ | References |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  | T | $T_{3}$ |  |  | ( MeV ) | $\left(m_{\pi^{\star}}\right)$ |  | (sec) | ( $1 / m_{\pi^{ \pm}}$) |  |  |  |
|  | $\begin{aligned} & \tilde{y}^{-} \\ & \Xi^{0} \end{aligned}$ | $\overline{\text { \% }}$ | ${ }^{\frac{1}{2}}$ | - $\begin{array}{r}\text { - } \\ \text { 2 } \\ \text { 2 } \\ \frac{1}{2} \\ \frac{1}{2} \\ \hline\end{array}$ | $\frac{1}{2}$ | $\begin{array}{r} -2 \\ 2 \\ -2 \end{array}$ | $\begin{gathered} 1320.8 \pm 0.4 \\ 1316 \end{gathered}$ | $\begin{aligned} & \hline 9.46 \\ & 9.43 \end{aligned}$ |  | $\begin{array}{\|c\|} \hline 1.4(+0.6 /-0.2) \\ \times 10^{-10} \\ 3.9(+1.4 /-0.9) \times 10^{-10} \end{array}$ | $\begin{aligned} & 3 \times 10^{13} \\ & 8 \times \ddot{10^{13}} . \end{aligned}$ | $\begin{aligned} & \Lambda \pi^{-} \\ & \Lambda \pi^{0} \end{aligned}$ | $\begin{aligned} & 100 \\ & 100 \end{aligned}$ | $1$ |
|  | $\begin{aligned} & \Sigma^{-} \\ & \Sigma^{0} \\ & \Sigma^{+} \end{aligned}$ | $\begin{aligned} & \bar{\Sigma}^{+} \\ & \bar{\Sigma}^{0} \\ & \bar{\Sigma}^{-} \end{aligned}$ | 1 | $\begin{gathered} -1 \\ 1 \\ 0 \\ 0 \\ 1 \\ -1 \end{gathered}$ | $\begin{aligned} & \frac{1}{2}+ \\ & \frac{1}{2}+ \\ & \frac{1}{2}+ \end{aligned}$ | $\begin{gathered} -1 \\ -1 \\ -1 \\ 1 \\ -1 \\ 1 \end{gathered}$ | $\begin{gathered} 1195.96 \pm 0.30 \\ 1191.5 \pm 0.5 \\ 1189.40 \pm 0.20 \end{gathered}$ | $\begin{aligned} & 8.57 \\ & 8.54 \\ & 8.52 \end{aligned}$ |  | $\begin{aligned} & (1.59 \pm 0.05) \times 10^{-10} \\ & 10^{-11}>\tau>10^{-22} \\ & (0.78 \pm 0.03) \times 10^{-10} \end{aligned}$ | $\left\lvert\, \begin{gathered} 3.4 \times 10^{13} \\ 10^{12}>r^{\prime}>10 \\ 1.65 \times 10^{13} \end{gathered}\right.$ | $\begin{gathered} n \pi^{-} \\ \Lambda \gamma \\ \underset{n \pi^{+}}{ } \end{gathered}$ | $\begin{gathered} 100 \\ 100 \\ 50.7 \pm 2.3 \\ 49.3 \pm 2.3 \end{gathered}$ | $\left\{\begin{array}{l} 3,4,19 \\ 3,5,19 \\ 3,4,19 \end{array}\right.$ |
|  | $\Lambda^{0}$ | $\bar{\Lambda}^{0}$ | 0 | $\begin{aligned} & 0 \\ & 0 \end{aligned}$ | ${ }^{\frac{1}{2}+}$ | $\left\|\begin{array}{c} -1 \\ 1 \end{array}\right\|$ | $\begin{aligned} & 1115.38 \pm 0.10 \\ & 1115.44 \pm 0.32 \end{aligned}$ | $\begin{aligned} & 7.991 \\ & 7.991 \end{aligned}$ | $-1.5 \pm 0.5$ | $\begin{gathered} (2.57 \pm 0.30) \times 10^{-10} \\ (1.9 \pm 1.0) \times 10^{-10} \end{gathered}$ | $\begin{gathered} 5.4 \times 10^{13} \\ 4 \times 10^{13} \end{gathered}$ | $p_{n \pi^{-}}{ }^{0}$ | $\left\|\begin{array}{l} 66(+4 /-3) \\ 34(+3 /-4) \end{array}\right\|$ | 6,20 |
| $\begin{aligned} & 0 \\ & 0 \\ & 0 \\ & 0 \\ & 0 \\ & 0 \\ & 7 \end{aligned}$ | $\begin{aligned} & n^{0} \\ & p^{+} \end{aligned}$ |  | ${ }^{\frac{1}{2}}$ | - $\begin{array}{r}\text { - } \\ \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \\ -\frac{1}{2} \\ \hline\end{array}$ | $\begin{aligned} & \frac{1}{2}+ \\ & \frac{1}{2}+ \end{aligned}$ | $\begin{aligned} & 0 \\ & 0 \\ & 0 \\ & 0 \end{aligned}$ | $\begin{aligned} 939.507 & \pm 0.01 \\ 938.213 & \pm 0.01 \end{aligned}$ | $\begin{aligned} & 6.731 \\ & 6.722 \end{aligned}$ | $\left\|\begin{array}{c} -1.9128 \\ 2.792816 \pm \\ 0.00034 \\ -1.8 \pm 1.2 \end{array}\right\|$ | $1013 \pm 26$ <br> $\infty$ | $2.15 \times 10^{26}$ <br> $\infty$ | $p e^{-\bar{\nu}_{e}}$ | 100 | $\begin{gathered} 7,8 \\ 7,15 \\ 10 \end{gathered}$ |
| $\begin{aligned} & \stackrel{u}{\tilde{Z}} \\ & \stackrel{\rightharpoonup}{0} \\ & \stackrel{y}{4} \end{aligned}$ | $K^{+}$ <br> $K^{0}$ | $\begin{aligned} & K^{-} \\ & \bar{K}^{0} \end{aligned}$ | ${ }^{\frac{1}{2}}$ |  | $0^{-}$ $0^{-}$ | $\left\|\begin{array}{c} 1 \\ \\ -1 \\ -1 \\ -1 \end{array}\right\|$ | $493.98 \pm 0.14$ $497.9 \pm 0.6$ | 3.539 $3.57$ | 0 $\underset{\hbar e_{i}, m_{K K}}{<0.04}$ | $\left\|\begin{array}{c} (1.227 \pm 0.008) \times 10^{-8} \\ \\ K_{1}^{0}(0.90 \pm 0.02) \times 10^{-10} \\ K_{2}^{0} 6.3(+1.6 /-1.0) \\ \times 10^{-8} \end{array}\right\|$ | $2.60 \times 10^{15}$ $\begin{aligned} & 1.9 \times 10^{13} \\ & 1.3 \times 10^{16} \end{aligned}$ |  | $64.2 \pm 1.3$ $18.6 \pm 0.9$ $4.8 \pm 0.6$ $5.0 \pm 0.5$ $5.7 \pm 0.3$ $1.7 \pm 0.2$ $69.4 \pm 1.0$ $30.6 \pm 1.0$ $8.7 \pm 2.3$ $38 \pm 7$ $28.3 \pm 5.9$ $25.0 \pm 5.9$ | 9 <br> 11 <br> 12 |
|  | $\begin{aligned} & \pi^{+} \\ & \pi^{0} \end{aligned}$ | $\frac{\pi^{-}}{\pi^{0}}$ | 1 | $\begin{gathered} 1 \\ -1 \\ 0 \end{gathered}$ | $\begin{aligned} & 0^{-} \\ & 0^{-} \end{aligned}$ | $\begin{aligned} & 0 \\ & 0 \\ & 0 \end{aligned}$ | $\begin{aligned} & 139.58 \pm 0.05 \\ & 134.97 \pm 0.05 \end{aligned}$ | $\begin{gathered} 1 \\ 0.967 \end{gathered}$ | $\begin{aligned} & 0 \\ & 0 \end{aligned}$ | $\begin{gathered} (2.547 \pm 0.027) \times 10^{-8} \\ (1.05 \pm 0.18) \times 10^{-16} \end{gathered}$ | $\begin{aligned} & 5.48 \times 10^{15} \\ & 2.23 \times 10^{7} \end{aligned}$ | $\begin{gathered} \mu^{+} \nu_{\mu} \\ 2 \gamma \\ 2 e^{+} e^{-} \end{gathered}$ | $\begin{gathered} 100 \\ 98.8 \\ 1.2 \end{gathered}$ | $\begin{gathered} 18 \\ 13,22 \end{gathered}$ |
|  | $\mu^{-}$ | $\mu^{+}$ |  |  | $\frac{1}{2}$ |  | 105.65 | $\left\|\begin{array}{c} 206.765 \\ \pm 0.002 \\ m_{e} \end{array}\right\|$ | $\begin{gathered} (1.001162 \pm \\ 0.000005) \\ e / 2 m_{\mu} \end{gathered}$ | $(2.210 \pm 0.002) \times 10^{-6}$ | $4.69 \times 10^{17}$ | $e^{-\nabla_{v} \nu_{\mu}}$ | 100 | 14 |
| $\begin{gathered} \stackrel{y}{0} \\ \stackrel{8}{0} \\ \stackrel{\rightharpoonup}{4} \\ \hline \end{gathered}$ | $e^{-}$ | ${ }_{-}^{+}$ |  |  | $\frac{1}{2}$ |  | $\begin{aligned} & 0.510976 \pm \\ & 0.000007 \end{aligned}$ | $1 m$ 。 | $\begin{gathered} (1.0011609 \pm \\ 0.0000024) \\ e / 2 m_{s} \end{gathered}$ | $\infty$ | $\infty$ |  |  | 7,15 |
|  | $\begin{aligned} & \nu_{\mu}{ }^{0} \\ & \nu_{\varepsilon}{ }^{0} \end{aligned}$ | $\tilde{\nu}_{\mu}$ <br> $\bar{\nu}_{e}$ |  |  | - <br> $-\frac{1}{2}$ <br> $+\frac{2}{2}$ <br> $-\frac{1}{2}$ <br> $+\frac{2}{2}$ |  | $\begin{aligned} & <2.5 \\ < & 0.00025 \end{aligned}$ | $\left\|\begin{array}{c} <5 m_{e} \\ <5 \\ \times 10^{-4} m_{e} \end{array}\right\|$ |  |  |  |  |  | $\begin{aligned} & 16,21 \\ & 17,21 \end{aligned}$ |
| $\begin{aligned} & 6 \\ & 0 \\ & 0 \\ & 0 \\ & \end{aligned}$ | $\gamma^{0}$ |  |  |  | 1 | 0 | 0 | 0 |  |  |  |  |  | 21. |

particles, except the muon and the neutrinos, for which it is expressed in $m_{e}$. The mass in the other unit gives only the significant figures.

The magnetic moment is expressed in proton magnetons for baryons, in muon magnetons for the muon, and in electron magnetons for the electron.
The mass and magnetic moments of antiparticles are included when they have been specifically measured; otherwise, a blank space is left. A blank space in the magnetic-moment column may also indicate that the value is not known.

## WIDTH AND LIFETIME

In Table I the mean life is given in two units, accurately in seconds and rounded off in $1 / m_{\pi \pm}$. The relation is

$$
1 / m_{\pi^{ \pm}}=4.7153 \times 10^{-24} \mathrm{sec}
$$

the latter time signifying the time required for light to travel the distance of a Compton wavelength of the $\pi^{ \pm}$meson. This distance equals

$$
1.4136 \times 10^{-13} \mathrm{~cm}
$$

In Table II the full width $\Gamma$ at half-maximum of the resonance is given in MeV , and the lifetime $\Gamma^{-1}$ in units of $1 / m_{\pi^{ \pm}}$, to allow comparison with the mean lives in Table I.

A blank space means that the quantity is not known. Widths and lifetimes of antiparticles, and, in Table II, charge multiplets, are not included unless they have been specifically measured.

## PRODUCTION PROPERTIES

In Table II one production reaction is given although others may also have been used. The laboratory momentum of the incident particle has been computed for that production reaction, at the threshold of resonance production.

A blank space in the $k_{\text {lab }}$ column signifies that the rest masses in the production reaction have not been computed because no threshold exists or because it is questionable which threshold is of interest.

A blank space in both production columns indicates that detailed information on the production of different charge states is not available.

## DECAY PROPERTIES

The most common decay modes are given if they have been observed. By "most common" we mean a branching ratio $\geqslant 1 \%$.

The decay modes of antiparticles are not listed because they are simply the antimodes of the particles, and the branching ratios are the same. A blank space in any of the decay-property columns signifies that the information is lacking.

## REFERENCES

All detailed information is collected in the reference list which does not claim to be complete. Use has been made of all literature available in Copenhagen by March 1963.

## ACKNOWLEDGMENTS

Table I is the third version of a table prepared in 1958 by Dr. Monica Hessler and Dr. Bertel Laurent and revised in 1960 by the present compiler, all at the Institute for Theoretical Physics, Stockholm, at that time. The reason why the first-mentioned authors do not appear as co-authors in the present version is that they must not be held responsible for any erroneous or incomplete information contained in the present Table I, the inheritance of which is hereby gratefully acknowledged.

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## REFEREINCES

1. $\mathrm{m}\left(\Xi^{-}\right)$is a weighted average of the following results in MeV:
$1320.4 \pm 2.2$, W. H. Barkas and A. H. Rosenfeld, University of California Radiation Laboratory Technical Report UCRL-8030 (unpublished), compilation of $12 \boldsymbol{\Xi}$ found before March 1958.
$1317.9 \pm 1.9$, W. B. Fowler, R. W. Birge, P. Eberhard, R. Ely, M. L. Good, W. M. Powell and H. K. Ticho, Phys. Rev. Letters 6, 134 (1961).
$1317.0 \pm 2.2$, V. A. Soloviev, Proceedings of the 1960 Annual International Conference on High Energy Physics at Rochester (University of Rochester, Rochester, 1960), p. 388.
$1321.0 \pm 0.5, \mathrm{~L}$. Bertanza, V. Brisson, P. L. Connolly, E. L. Hart, I. S. Mittra, G. C. Moneti, R. R. Rau, N. P. Samios, I. O. Skillicorn, S. S. Yamamoto, M. Goldberg, L. Gray, J. Leitner, S. Lichtman, and J. Westgard, Phys. Rev. Letters 9, 229 (1962).
$1321.2 \pm 2.4$, a group from CERN, Ecole Polytechnique, and CEN, Saclay, Phys. Rev. Letters 8, 257 (1962). A measurement of $\vec{E}^{+}$.
$1322.0 \pm 1.3$, H. N. Brown, B. B. Culwick, W. B. Fowler, M. Gailloud, T. E. Kalogeropoulos, J. K. Kopp, R. M. Lea, R. I. Louttit, T. W. Morris, R. P. Shutt, A. M. Thorndike, M. S. Webster, C. Baltay, E. C. Fowler, J. Sandweiss, J. R. Sanford, and H. D. Taft, Phys. Rev. Letters 8, 255 (1962). A measurement of ${ }_{\boldsymbol{\Xi}}+$.
$\tau\left(\Xi^{-}\right)$is a central value of the following results in $10^{-10} \mathrm{sec}:$ $1.28(+0.41 /-0.25)$, Fowler, et al. (cf. above),
$1.16(+0.26 /-0.17)$, Bertanza, et al. (cf. above).
$1.91(+0.35 /-0.25)$, L. Jauneau, D. Morellet, U. NguyenKhac, P. Petiau, A. Rousset, H. Bingham, D. C. Cundy, W. Koch, B. Ronne, H. Sletten, F. W. Bullock, A. K. Common, M. J. Esten, C. Henderson, F. R. Stannard, J. M. Scarr, J. Sparrow, A. G. Wilson, A. Halsteinslid, and R. Möllerud, Phys. Letters 4, 49 (1963)
$J=1 / 2$ is compatible with the data of Bertanza, et al.
2. $m\left(\Xi^{0}\right)$ is privately communicated by A. H. Rosenfeld, who refers to F. T. Solmitz' talk at Stanford American Physical Society meeting, Dec. 1962. $\tau\left(\widetilde{\Xi}^{0}\right)$ is from Jauneau et al. (cf. reference 1).
3. The $\Sigma$ mass has been taken from the compilation of W. H. Barkas and A. H. Rosenfeld, Proceedings of the 1960 Annual International Conference on High Energy Physics at Rochester (University of Rochester, Rochester, 1960), p. 877.
4. $\tau\left(\Sigma^{+}\right)$and $\tau\left(\Sigma^{-}\right)$are weighted averages of the following results in $10^{-10} \mathrm{sec}$ :
$\tau\left(\Sigma^{-}\right)=1.61_{-0.09}^{+0.0}, \tau\left(\Sigma^{+}\right)=0.81_{-0.05}^{+0.06}$, Barkas and Rosenfeld compilation (cf. reference 3 ).
$\tau\left(\Sigma^{-}\right)=1.58 \pm 0.06, \tau\left(\Sigma^{+}\right)=0.765 \pm 0.04$, W. E. Humphrey and R. R. Ross, Phys. Rev. 127, 1305 (1962).
The $\Sigma^{+}$branching ratio for $\mathrm{p} \pi^{0}$ decay is a weighted average of the following results: $(51.0 \pm 2.4) \%$, Humphrey and Ross; $(48 \pm 7) \%$, P. Granet, Compt. Rend. 255, 282 (1962).
5. The upper limit of $\tau\left(\Sigma^{0}\right)$ is from L. W. Alvarez, H. Bradner, P. Falk-Variant, J. D. Gow, A. H. Rosenfeld, F. T. Solmitz, and R. Tripp, University of California Radiation Laboratory Technical Report UCRL-3775 (unpublished); the lower limit, from J. Dreitlein and B. W. Lee, Phys. Rev. 124, 1274 (1961), who also give a theoretical estimate, $\tau\left(\Sigma^{0}\right)=$ $1.1 \times 10^{-19}$.
6. $\mathrm{m}(\Lambda)$ is a weighted average of the following results in MeV:
$1115.36 \pm 0.14$, compilation of Barkas and Rosenfeld (cf. reference 3).
$1115.46 \pm 0.15$, B. Bhowmik, D. P. Goyal and N. K. Yamdagni, Nuovo Cimento 22, 296 (1961).

Table IIa. Baryonic Resonant States, March 1963

| $\begin{gathered} \frac{4}{5} \\ \text { E } \\ \hline \mathbf{E} \end{gathered}$ | $\begin{aligned} & \text { ত̈ } \\ & \text { 俞 } \\ & \text { N } \end{aligned}$ |  |  |  | $S$ | Mass |  | $\begin{gathered} \text { Full } \\ \text { width } \Gamma \\ (\mathrm{MeV}) \end{gathered}$ | Lifetime $\mathrm{T}^{-1}$ ( $1 / m_{\pi^{ \pm}}$) | Production |  | Decay |  |  | References |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |  |  |  |  |  | Modes | Branchin |  |  |
|  |  |  |  |  |  | ( MeV ) | ( $m_{\pi^{ \pm}}$) |  |  | Process | ( MeV ) |  | ratio (\%) | ( MeV ) |  |
|  | $Y_{05}^{*}$ | 0 | 0 | $\frac{5}{2}$ or $\frac{3}{2}$ | -1 | 1815 | 13.0 |  | 120 | 1.16 | $K^{-p}$ | 1050 | $K p$ others |  | 383 | 23 |
|  | $Y^{* *}$ | 0 |  |  | $-1$ | $1770 \pm 100$ | 12.7 |  |  | $\pi^{-} p$ | 2260 | $\begin{aligned} & Y_{13}^{*+} \pi^{-} \\ & \Lambda \pi^{+} \pi^{-} \end{aligned}$ |  | 245 376 | 55 |
|  | $Y^{*}$ | 0 |  |  | -1 | 1715 | 12.2 |  |  | $\pi \sim$ | 2185 | $K^{0} n$ | 100 | 297 | 24 |
|  | $Y_{0}^{* *}$ | 0 | 0 |  | $-1$ | 1680 | 12.0 | $<20$ | $>7$ | $K^{-} p$ | 760 | $\Lambda \eta$ |  | 16 | 48 |
| ¢ | $Y_{13}^{* *}$ | + | 1 | $\frac{3}{2}$ | -1 | $1660 \pm 10$ | 11.9 | $40 \pm 10$ | 3.5 | $K-p$ | 715 | $\begin{gathered} \bar{K}^{0} p \\ (\Sigma \pi)^{+} \\ \Lambda \pi^{+} \\ \Lambda \pi^{+} \pi^{0} \\ \Sigma^{\mp} \pi^{ \pm} \pi^{+} \end{gathered}$ | $\begin{array}{r} \sim 10 \\ 30 \\ 25 \\ 20 \\ 15 \end{array}$ | $\begin{aligned} & 224 \\ & 333 \\ & 405 \\ & 270 \\ & 188 \end{aligned}$ | 56 |
| . | $Y_{2}^{*}$ |  | 1 or 2 |  | -1 | $1550 \pm 20$ | 11.1 | 125 | 1.75 | $\pi^{-1} p$ | 1770 | $\Sigma \pi$ | 100 | 227-236 | 25 |
|  | $\begin{aligned} & \Xi_{1}^{*} \\ & z_{1}^{*} \end{aligned}$ | $\begin{aligned} & - \\ & 0 \end{aligned}$ | $\frac{1}{2}$ | $\frac{3}{2}+$ | $\begin{aligned} & -2 \\ & -2 \end{aligned}$ | $1533 \pm 3$ | 10.98 | $\leq 7$ | $\geq 20$ | $K-p$ | $\begin{aligned} & 1512 \\ & 1521 \end{aligned}$ |  | $\begin{array}{r} 40 \\ 60 \\ 100 \end{array}$ | 78 82 73 | $\begin{aligned} & 26 \\ & 26 \end{aligned}$ |
|  | $Y_{03}^{*}$ | 0 | 0 | $\frac{3}{2} \quad-$ | -1 | $1520 \pm 3$ | 10.89 | 16 | 8.7 | $K-p$ | 395 | $\begin{gathered} (\bar{K} N)^{0} \\ (\Sigma \pi)^{0} \\ \Lambda \pi^{+} \pi^{-} \end{gathered}$ | 33 56 11 | $\begin{gathered} 82-88 \\ 184-193 \\ 126 \end{gathered}$ | 27 |
|  | $Y_{0}^{*}$ | 0 | 0 |  | -1 | $\begin{aligned} & 1404.7 \pm \\ & 0.4 \end{aligned}$ | 10.1 | $<1.4$ | $>100$ | $K \sim p$ | 445 | $\begin{gathered} (\Sigma \pi)^{0} \\ \Lambda \pi \pi \end{gathered}$ |  | $\begin{aligned} & 69-78 \\ & 10-20 \end{aligned}$ | 28 |
|  | $\underline{Y}$ | - 0 + | 1 | $\frac{1}{2}$ or ${ }^{3}+$ | $\begin{aligned} & -1 \\ & -1 \\ & -1 \end{aligned}$ | $1385 \pm 5$ | 9.92 | $50 \pm 10$ | 2.8 | $K-p$ | $\begin{aligned} & 408 \\ & 395 \\ & 408 \end{aligned}$ | $\begin{gathered} (\Sigma \pi)^{-} \\ A \pi^{-} \\ (\Sigma \pi)^{0} \\ \Lambda \pi^{3} \\ (\Sigma \pi)^{+} \\ \Lambda \pi^{+} \end{gathered}$ | $\begin{array}{r} 1( \pm 3) \\ 99( \pm 3) \\ \\ 1( \pm 3) \\ 99( \pm 3) \end{array}$ | $\begin{gathered} 54 \\ 130 \\ 49-58 \\ 135 \\ 53-61 \\ 130 \end{gathered}$ | 29 <br> 29 <br> 29 |
|  | $N_{3}^{*}$ |  | $\frac{3}{2}$ |  | 0 | $2360 \pm 25$ | 16.9 | $200 \pm 25$ | 0.7 | $\pi^{+} p$ | 2510 | $\pi N$ others |  | 1280 | 60 |
|  | $N_{i}^{*}$ |  | $\frac{1}{2}$ |  | 0 | $2190 \pm 25$ | 15.7 | $200 \pm 20$ | 0.7 | $\pi^{-} p$ | 2080 | $\pi N$ <br> others |  | 1110 | 60 |
|  | $Z_{3}^{*}$ | 0 | $\frac{3}{2}$ | $\geqslant \frac{3}{2}$ | 0 | $1920 \pm 20$ | 13.8 | 15 | 9 | $\pi^{-}(\mathrm{A})$ |  | $\begin{gathered} K^{0} \Lambda \\ (K \Sigma)^{0} \end{gathered}$ |  | $\begin{aligned} & 307 \\ & 231 \end{aligned}$ | $\begin{aligned} & 30 \\ & 57 \end{aligned}$ |
|  | $N_{37}^{*}$ |  | $\frac{3}{2}$ | $\frac{7}{2}$ | 0 | 1900 | 13.6 | 200 | 0.7 | $\pi \sim(A)$. | 1440 | $\frac{\pi N}{K \Sigma}$ | $\begin{aligned} & 30 \\ & <4 \end{aligned}$ | $\begin{aligned} & 820 \\ & 215 \end{aligned}$ | 31 |
| \% | $N_{1 i}^{*}$ |  | $\frac{1}{2}$ | $\frac{1}{2}+$ | 0 | 1690 | 12.1 |  |  | $\pi^{-} p$ | 1030 | $\pi N$ | 100 | 612 | 32 |
| $\begin{aligned} & 2 \\ & 0 \\ & \ddot{0} \\ & 0 \\ & 0 \\ & \hline \end{aligned}$ | $N_{15}^{*}$ |  | $\frac{1}{2}$ | $\frac{5}{2}+$ | 0 | $1683 \pm 5$ | 12.06 | 80 | 1.7 | $\pi \sim$ | 1020 | $\pi N$ <br> K others | $\begin{gathered} 80 \\ <2 \\ >18 \\ \hline \end{gathered}$ | $\begin{array}{r} 605 \\ 74 \end{array}$ | 31 |
| 学 | $Z_{1}^{*}$ | 0 | . ${ }^{\frac{1}{2}}$ | $\geqslant \frac{1}{2}$ | 0 | $1650 \pm 20$ | 11.8 | $<7$ | $>20$ | $\pi^{-}(A)$ |  | $K^{0} \Lambda$ | 100 | 38 | 30 |
|  | $N_{13}^{*}$ |  | $\frac{1}{2}$ | $\frac{3}{2} \quad-$ | 0 | $1517 \pm 3$ | 10.87 | 60 | 2.3 | $\pi^{-} p$ | 731 | $\pi N$ <br> others |  | 439 | 31 |
|  | $\begin{aligned} & N_{33}^{*} \\ & N_{33}^{*} \\ & N_{33}^{*} \\ & N_{33}^{*} \end{aligned}$ | $\begin{gathered} - \\ 0 \\ + \\ ++ \end{gathered}$ | $\frac{3}{2}$ | $\frac{3}{2}+$ | $\begin{aligned} & 0 \\ & 0 \\ & 0 \\ & 0 \end{aligned}$ | 1237 | 8.86 | $90 \pm 20$ | 1.6 | $\pi \mathrm{V}$ | 303 | $\begin{aligned} & \pi^{-n} n \\ & \pi^{-p} p \\ & \pi^{0} n \\ & \pi^{0} p \\ & \pi^{+} n \\ & \pi^{+} p \end{aligned}$ | $\begin{aligned} & 100 \\ & 100 \end{aligned}$ | $\begin{aligned} & 158 \\ & 159 \\ & 163 \\ & 164 \\ & 158 \\ & 159 \end{aligned}$ | $\begin{aligned} & 31 \\ & 31 . \\ & 31 \\ & 31 \end{aligned}$ |

$1115.25 \pm 0.36$, R. Armenteros, E. Fett, B. French, L. Montanet, V. Nikitin, M. Szeptycka, Ch. Peyrou, R. Böck, A. Shapira, J. Badier, L. Blaskovicz, B. Equer, B. Gregory, F. Muller, S. J. Goldsack, D. H. Miller, C. C. Butler, B. Tallini, J. Kinson, L. Riddiford, A. Leveque, J. Meyer, A. Verglas, and S. Zylberach, Proceedings of the 1962 Annual International Conference on High-Energy Physics at Geneva (CERN, Geneva, Switzerland, 1962), p. 236.
$1115.04 \pm 0.41$, C. Baltay, E. C. Fowler, J. Sandweiss, J. R. Sanford, H. D. Taft, B. B. Culwick, W. F. Fowler, J. K. Kopp, R. I. Louttit, R. P. Shutt, A. M. Thorndike, and M. S. Webster, ibid., p. 233.
$\mu(\Lambda)$ is the value of R. L. Cool, E. W. Jenkins, T. F. Kycia, D. A. Hill, L. Marshall, and R. A. Schluter, Phys. Rev. 127, 2223 (1962). A value of $0.0 \pm 0.6$ nucleon magnetons, given by W. Kernan, T. B. Novey, S. D. Warshaw, and A. Wattenburg, Phys. Rev. 129, 870 (1963), has not been used.
$\tau(\Lambda)$ is a central value from the ideogram of F.S. Crawford, Proceedings of the 1962 Annual International Conference on High-EnergyPhysicsatGeneva (CERN,Geneva, Switzerland,1962),p.827.
$m(\bar{\Lambda})$ is a weighted average of the following results in MeV :
$1115.40 \pm 0.39$, Baltay et al. (cf. above).
$1115.52 \pm 0.55$, Armenteros et al. (cf. above).
$r(\bar{\Lambda})$ is from Baltay et al.
The $\Lambda \rightarrow \mathrm{p} \pi^{-}$branching ratio is a central value of the following results in percent:
$64.3 \pm 1.6$, Humphrey and Ross (cf. reference 4),
$64.5 \pm 2.2$, weighted average based on earlier measurements, quoted by F. S. Crawford (cf. above),
$68.5 \pm 1.7$, J. A. Anderson, F. S. and B. B. Crawford, R. L. Golden, L. J. Lloyd, G. W. Meisner and L. Price (to be published), quoted by F. S. Crawford (cf. above).
7. The nucleon and electron masses and the neutron magnetic moment are from compilations by Cohen, Crowe,

Table IIb. Mesonic Resonant States, March 1963.

\begin{tabular}{|c|c|c|c|c|c|c|c|c|c|c|c|c|c|c|}
\hline \multirow[b]{3}{*}{} \& \multirow[b]{3}{*}{} \& \multirow[b]{3}{*}{$$
\begin{aligned}
& \text { 荷 } \\
& \text { 品 }
\end{aligned}
$$} \& \multirow[t]{3}{*}{} \& \multirow{3}{*}{S} \& \multicolumn{2}{|l|}{\multirow[b]{2}{*}{Mass}} \& \multirow[b]{3}{*}{$$
\begin{aligned}
& \text { Full } \\
& \text { width } \mathrm{T} \\
& (\mathrm{MeV})
\end{aligned}
$$} \& \multirow[b]{3}{*}{$$
\begin{gathered}
\text { Life- } \\
\text { time } \Gamma^{-1} \\
\left(1 / m_{\pi^{ \pm}}\right)
\end{gathered}
$$} \& \multicolumn{2}{|l|}{Production} \& \multicolumn{3}{|c|}{Decay} \& \multirow[b]{3}{*}{References} <br>
\hline \& \& \& \& \& \& \& \& \& \multirow[b]{2}{*}{Process} \& \multirow[t]{2}{*}{$$
\begin{gathered}
\mathcal{K}_{\text {lab }} \\
(\mathrm{MeV})
\end{gathered}
$$} \& \multirow[t]{2}{*}{Modes} \& \multirow[t]{2}{*}{Branching
ratio (\%)} \& \multirow[t]{2}{*}{$$
\begin{gathered}
Q \\
(\mathrm{MeV})
\end{gathered}
$$} \& <br>
\hline \& \& \& \& \& $(\mathrm{MeV})$ \& $\left(m_{\pi^{ \pm}}\right)$ \& \& \& \& \& \& \& \& <br>
\hline $\bar{K}_{3}^{*}$
$K_{3}^{*}$ \& + \& $\geq \frac{3}{2}$ \& \& 1
1 \& $1630 \pm 100$ \& 11.7 \& \& \& $\pi^{-} p$ \& 3534 \& $$
\begin{gathered}
\left(K_{1}^{*} \pi \pi\right)^{-} \\
\left(K_{2}{ }^{2}\right)- \\
\left(K_{\mu \rho}\right)- \\
\text { others } \\
\text { same, charge }+
\end{gathered}
$$ \& \& $$
\begin{aligned}
& 470 \\
& \leqslant 100 \\
& 225 \\
& \text { same }
\end{aligned}
$$ \& 55
55 <br>
\hline $\chi_{2}$ \& 0 \& \& \& 0 \& $1340 \pm 70$ \& 9.6 \& \& \& $\pi^{-p}$ \& 2287 \& $$
\begin{aligned}
& (\rho \pi \pi)^{0} \\
& \text { others }
\end{aligned}
$$ \& \& 290 \& 55 <br>
\hline $\kappa_{3}$ \& 0 \& \& \& 0 \& $1275 \pm 25$ \& 9.1 \& \& \& $\pi \sim$ \& 2125 \& ${ }^{K^{0} \bar{K}^{0}}{ }^{+} K^{-}$ \& \& 279
287 \& 24 <br>
\hline $K^{* *}$ \& \& \& \& 1 \& 1260 \& 9.0 \& \& \& $\pi-N$ \& \& $K(n \pi)$ \& \& \& 55 <br>
\hline $f$ \& 0 \& 0 \& $2++$ \& 0 \& $1253 \pm 20$ \& 9.0 \& $100 \pm 50$ \& 1.4 \& $\pi^{-p}$ \& 2070 \& $\pi^{-} \pi^{+}$ \& 100 \& 970 \& 33 <br>
\hline $$
\begin{aligned}
& K_{5}^{*} \\
& K_{5}^{*}
\end{aligned}
$$ \& ++ \& $\frac{5}{2}$ \& \& ${ }_{1}^{1}$ \& $1150 \pm 50$ \& 8.2 \& \& \& $\pi^{-p}$ \& 2250 \& $$
\begin{aligned}
& K^{0} \pi^{+} \pi^{4} \\
& K^{0} \pi^{-} \pi^{-}
\end{aligned}
$$ \& \& 373
373 \& $$
\begin{aligned}
& 55 \\
& 55
\end{aligned}
$$ <br>
\hline $\chi$
$\chi_{1}$

1 \& - \& 1 \& \& 0 \& $$
\begin{aligned}
& 1050 \\
& 1040
\end{aligned}
$$ \& \[

$$
\begin{aligned}
& 7.5 \\
& 7.4
\end{aligned}
$$

\] \& \& \& $\pi^{-p}$ \& 1620 \& \[

$$
\begin{aligned}
& \pi^{-\pi} \pi^{-} \pi^{+} \pi^{0} \\
& \pi^{+} \pi^{-}(\pi \pi)^{0}
\end{aligned}
$$

\] \& \& \[

$$
\begin{gathered}
496 \\
481-491
\end{gathered}
$$

\] \& \[

$$
\begin{aligned}
& 51 \\
& 53
\end{aligned}
$$
\] <br>

\hline $\kappa_{2}$ \& 0 \& 0 \& even ++ \& 0 \& $1040 \pm 40$ \& 7.4 \& \& \& $K^{-p}$ \& 1780 \& $$
\begin{gathered}
\mathrm{K}_{1}^{0} K_{1}^{\circ} \\
\text { even number } \pi \text { 's }
\end{gathered}
$$ \& \& 44 \& 58 <br>

\hline $\kappa_{1}$ \& 0 \& 0 \& odd -- \& 0 \& 1020 \& 7.3 \& $<3$ \& $>47$ \& $K^{-} p$ \& 1760 \& $$
\begin{gathered}
K_{1}^{0} K_{2}^{0} \\
\text { odd number } \pi \text { 's }
\end{gathered}
$$ \& \& 24 \& 59 <br>

\hline \[
$$
\begin{aligned}
& \psi_{5} \\
& \psi_{5} \\
& \psi_{5}
\end{aligned}
$$

\] \& | - |
| :--- |
| - |
| + |
| + | \& 2 \& \& 0

0
0 \& 990 \& 7.2 \& \& \& $\pi \sim$ \& 1490 \& $\pi^{--} \pi^{-}$
$\pi^{-} \pi^{+}$

$\pi^{+} \pi^{+}$ \& $$
\begin{aligned}
& 100 \\
& 100
\end{aligned}
$$ \& 711

711

711 \& | 52 |
| :--- |
| 52 |
| 52 | <br>

\hline $K_{1}^{*}$ \& - \& $\frac{1}{2}$ \& 1 - \& -1 \& $888 \pm 3$ \& 6.4 \& $50 \pm 10$ \& 2.8 \& $K^{\prime} p$ \& 1074 \& $\bar{K}^{0} \pi^{-}$ \& $60 \pm 16$ \& 252 \& 34 <br>
\hline $\bar{K}_{1}^{*}$ \& 0 \& \& \& -1 \& \& \& \& \& K-p \& 1078 \& ${ }_{\text {K- }}{ }^{-\pi^{+}}$ \& \& 256 \& 34 <br>
\hline $K_{1}^{*}$ \& $+$ \& \& \& 1 \& \& \& \& \& $\pi^{-} p$ \& 1834 \& $K^{0} \pi^{+}$ \& 67 \& 258
252 \& 34 <br>
\hline \& 0 \& \& \& \& \& \& \& \& \& \& ${ }^{K^{+} \pi^{0}}{ }^{+} \pi^{-}$ \& 33 \& 261 \& <br>
\hline \& \& \& \& \& \& \& \& \& \& \& $K^{0} \pi^{0}$ \& \& 257 \& <br>
\hline $\varphi{ }^{3}$ \& 0 \& \& \& 0 \& $885 \pm 10$ \& 6.3 \& \& \& $\pi^{-} p$ \& 1284 \& $\pi^{+} \pi^{-}$ \& \& 606 \& 36 <br>

\hline $\omega$. \& $\bigcirc$ \& 0 \& 1 - - \& 0 \& $781.1 \pm 0.8$ \& 5.6 \& $<12$ \& > 12 \& $\bar{p} p$ \& \& \[
$$
\begin{gathered}
\text { neutr. } \\
\pi^{+} \pi^{-} \pi^{0} \\
\pi^{+} \pi^{0} \gamma^{0} \\
\pi^{+}+\pi^{-} \pi^{0} \pi^{0} \\
\pi^{+} \pi^{-} \pi^{+} \pi^{-} \\
\pi^{+} \pi^{-}
\end{gathered}
$$

\] \& \[

$$
\begin{gathered}
0.12 \pm 0.03 \\
<2 \\
<12 \\
<5 \\
4
\end{gathered}
$$

\] \& \[

$$
\begin{aligned}
& 373 \\
& 503 \\
& 232 \\
& 223 \\
& 503
\end{aligned}
$$

\] \& \[

$$
\begin{aligned}
& 37 \\
& 53 \\
& 53
\end{aligned}
$$
\] <br>

\hline $\rho$
$\rho$ \& - \& 1 \& $1-+$ \& 0
0 \& $757 \pm 5$
$751 \pm 6$ \& 5.4
5.4 \& $120 \pm 10$
$110 \pm 10$ \& 1.2

1.3 \& $\pi^{-p}$ \& 1029 \& \[
$$
\begin{gathered}
\pi^{-}-\pi^{0} \\
\pi^{-} \pi^{0} \pi^{0} \\
\pi^{-}-\pi^{0} \pi^{0} \pi^{0} \\
\pi^{-} \pi^{+} \pi^{-} \pi^{0}
\end{gathered}
$$

\] \& \[

$$
\begin{aligned}
& >91 \\
& <3 \\
& \leq 4 \\
& <2
\end{aligned}
$$

\] \& \[

$$
\begin{aligned}
& 475 \\
& 340 \\
& 205 \\
& 196
\end{aligned}
$$

\] \& \[

$$
\begin{aligned}
& 38,33,39 \\
& 53 \\
& 54 \\
& 53
\end{aligned}
$$
\] <br>

\hline $\rho$ \& 0 \& \& \& 0 \& $751 \pm 6$ \& 5.4 \& $110 \pm 10$ \& 1.3 \& $\pi N$ \& 1029 \& \[
$$
\begin{gathered}
\pi^{-} \pi^{+} \\
\text {neutrals }
\end{gathered}
$$

\] \& \[

$$
\begin{gathered}
94(+6 /-40) \\
6(+40 /-6)
\end{gathered}
$$
\] \& 470 \& 40, 33, 36, 39 <br>

\hline $\rho_{2}$ \& 0 \& \& \& 0 \& 780 \& 5.6 \& 60 \& 2.3 \& $\pi N$ \& 1085 \& $\pi^{+-\pi^{--} \pi^{+} \pi^{-}} \pi^{-\pi^{+}}$ \& <2 \& \[
$$
\begin{aligned}
& 191 \\
& 500
\end{aligned}
$$

\] \& \[

$$
\begin{aligned}
& 53 \\
& 41
\end{aligned}
$$
\] <br>

\hline $\rho_{1}$ \& 0 \& \& \& 0 \& 720 \& 5.2 \& 20 \& 7 \& $\pi N$ \& 975 \& $\pi^{-\pi^{+}}$ \& \& 440 \& 40 <br>
\hline $\rho$ \& + \& \& \& 0 \& \& \& \& \& $\pi^{+} p$ \& 1066 \& \& \& 495 \& 39 <br>

\hline $$
\begin{aligned}
& \psi_{1} \\
& \psi_{4} \\
& \psi_{4}
\end{aligned}
$$ \& \[

$$
\begin{gathered}
- \\
0 \\
++
\end{gathered}
$$
\] \& 2 \& \& 0

0

0 \& 760 \& 5,4 \& \& \& $\pi^{-} p$ \& \[
$$
\begin{aligned}
& 1310 \\
& 1055 \\
& 1590
\end{aligned}
$$

\] \& \[

$$
\begin{aligned}
& \pi^{-} \pi^{-} \\
& \pi^{-} \pi^{+} \\
& \pi^{+} \pi^{+}
\end{aligned}
$$

\] \& \[

$$
\begin{aligned}
& 100 \\
& 100
\end{aligned}
$$

\] \& \[

$$
\begin{aligned}
& 481 \\
& 481 \\
& 481
\end{aligned}
$$

\] \& \[

$$
\begin{aligned}
& 52 \\
& 52 \\
& 52
\end{aligned}
$$
\] <br>

\hline \& 0 \& $\frac{1}{2}$ \& $\geqslant 1$ \& 1 \& $730 \pm 10$ \& 5.2 \& $\leq 20$ \& $\geqslant 7$ \& $\pi \sim p$ \& 1485 \& $\stackrel{K^{+} \pi^{-}}{K^{0} \pi^{0}}$ \& \& 96
97 \& 50, 56 <br>
\hline $K_{1}^{* *}$ \& + \& \& \& 1 \& \& \& \& \& \& \& $(K \pi)^{+}$ \& \& 92-101 \& 50, 56 <br>
\hline $\delta$ \& $\bigcirc$ \& 1 or 2 \& \& 0
0 \& $645 \pm 25$ \& 4.5 \& \& \& $\pi^{-} p$ \& 810 \&  \& \& 350
345 \& 43
43 <br>
\hline $\delta$ \& + \& \& \& 0 \& \& \& \& \& \& \& $\pi^{+} \pi^{0}$ \& \& 350 \& 43 <br>
\hline ${ }_{\alpha}^{\alpha}$ \& 0
+

+ \& 1 or 2 \& \& 0

0 \& 625 \& 4.5 \& <80 \& $>1.7$ \& $p p$ \& \& \[
\frac{\pi^{+} \pi^{-} \pi^{0}}{\pi^{+} \pi^{+} \pi^{-}}

\] \& \& 225 \& \[

$$
\begin{aligned}
& 42 \\
& 42
\end{aligned}
$$
\] <br>

\hline \[
$$
\begin{aligned}
& \psi_{3} \\
& \psi_{3} \\
& \psi_{3}
\end{aligned}
$$

\] \& | - |
| :--- |
|  |
| ++ |
| + | \& 2 \& 0 or 2 \& 0

0
0 \& $605 \pm 25$

580 \& $$
\begin{aligned}
& 4.3 \\
& 4.2
\end{aligned}
$$ \& 75 \& 1.9 \& $\pi-p$ \& \[

$$
\begin{array}{r}
1025 \\
733 \\
1235
\end{array}
$$

\] \& \[

$$
\begin{aligned}
& \pi^{-} \pi^{-} \\
& \pi^{-} \pi^{+} \\
& \pi^{+} \pi^{+}
\end{aligned}
$$

\] \& \[

$$
\begin{aligned}
& 100 \\
& 100
\end{aligned}
$$

\] \& \[

$$
\begin{array}{r}
326 \\
.301 \\
326
\end{array}
$$

\] \& \[

$$
\begin{aligned}
& 45,52 \\
& 52 \\
& 52,45
\end{aligned}
$$
\] <br>

\hline $\zeta$

$\zeta$ \& | - |
| :--- |
|  |
| + |
| + | \& 1 \& \& 0

0
0 \& $564 \pm 9$

$541 \pm 18$ \& \[
$$
\begin{aligned}
& 4.0 \\
& 3.9
\end{aligned}
$$

\] \& $<43$ \& >3.2 \& \[

$$
\begin{aligned}
& \pi^{-} p \\
& \pi^{-p} p \\
& \pi^{+} p
\end{aligned}
$$

\] \& \[

$$
\begin{aligned}
& 707 \\
& 672
\end{aligned}
$$

\] \& \[

$$
\begin{aligned}
& \pi^{-\pi} \pi^{0} \\
& \pi^{+} \pi^{-} \\
& \pi^{+} \pi^{0}
\end{aligned}
$$

\] \& \& \[

$$
\begin{aligned}
& 289 \\
& 262 \\
& 289
\end{aligned}
$$

\] \& \[

$$
\begin{aligned}
& 44 \\
& 44 \\
& 44
\end{aligned}
$$
\] <br>

\hline
\end{tabular}

| $\begin{aligned} & \stackrel{\rightharpoonup}{0} \\ & \stackrel{\text { a }}{6} \end{aligned}$ |  |  |  | S | Mass |  | $\begin{gathered} \text { Full } \\ \text { width } \Gamma \\ (\mathrm{MeV}) \end{gathered}$ | $\left.\begin{array}{\|c\|} \text { Life-- } \\ \text { time } \Gamma^{-1} \\ \left(1 / m_{\pi^{ \pm}}\right) \end{array} \right\rvert\,$ | Production |  | Decay |  |  | References |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |  | Process |  | $\begin{gathered} k_{i_{a b}} \\ (\mathrm{MeV}) \\ \hline \end{gathered}$ |  | Branching ratio (\%) | $\begin{gathered} Q \\ (\mathrm{MeV}) \end{gathered}$ |  |
|  |  |  |  |  | ( MeV ) | $\left(m_{\pi^{ \pm}}\right)$ |  |  |  |  |  |  |  |
| $\eta$ | 0 | 0 | $0-+$ | 0 | $548.5 \pm 0.6$ | 3.93 | $\leq 7$ | $\geqslant 20$ | $\pi p$ | 685 | $\begin{gathered} \pi^{+} \pi^{-} \pi^{0} \\ \pi^{+} \pi^{-} \gamma \\ 3 \pi^{0} \\ \pi^{0} \gamma \\ 2 \gamma \\ \text { others } \end{gathered}$ | $\begin{aligned} 25 & \pm 10 \\ 7 & \pm 2 \\ 68 & \pm 10 \end{aligned}$ | $\begin{aligned} & 135 \\ & 270 \\ & 144 \\ & 414 \\ & 549 \end{aligned}$ | 46 |
| $\varphi_{2}$ | 0 | 0 |  | 0 | $520 \pm 20$ | 3.7 | $70 \pm 30$ | 2.0 | $\pi^{-p}$ | 639 | $\pi^{+} \pi^{-}$ |  | 240 | 47 - |
| $\begin{aligned} & \psi_{2} \\ & \psi_{2} \\ & \psi_{2} \end{aligned}$ | $\begin{aligned} & - \\ & 0 \\ & ++ \end{aligned}$ | 2 |  | 0 <br> 0 <br> 0 | $\begin{aligned} & 440 \\ & 420-440 \\ & 440 \end{aligned}$ | $\begin{aligned} & 3.1 \\ & 3.1 \\ & 3.1 \end{aligned}$ |  |  | $\pi^{-} p$ | $\begin{aligned} & 735 \\ & 515 \\ & 975 \end{aligned}$ | $\begin{aligned} & \pi^{-} \pi^{-} \\ & \pi^{-} \pi^{+} \\ & \pi^{+} \pi^{+} \end{aligned}$ | $\begin{aligned} & 100 \\ & 100 \end{aligned}$ | $\begin{aligned} & 160 \\ & 160 \\ & 160 \end{aligned}$ | $\begin{aligned} & 52 \\ & 45,52 \\ & 52 \end{aligned}$ |
| $\varphi_{1}$ | 0 | 0 |  | 0 | $395 \pm 10$ | 2.8 | $50 \pm 20$ | 2.8 | $\pi^{-p}$ | 446 | $\pi^{+} \pi^{-}$ |  | 115 | 47 |
| $\begin{aligned} & \psi_{1} \\ & \psi_{1} \\ & \psi_{1} \\ & \hline \end{aligned}$ | $\begin{aligned} & - \\ & 0 \\ & ++ \end{aligned}$ | 2 |  | 0 0 0 | $\begin{aligned} & 330 \\ & 330 \\ & 330 \\ & \hline \end{aligned}$ | $\begin{aligned} & 2.4 \\ & 2.4 \\ & 2.4 \end{aligned}$ |  |  | $\pi^{-} p$ | $\begin{aligned} & 557 \\ & 346 \\ & 790 \end{aligned}$ | $\begin{aligned} & \pi^{-} \pi^{-} \\ & \pi^{-} \pi^{+} \\ & \pi^{+} \pi^{+} \end{aligned}$ | $\begin{aligned} & 100 \\ & 100 \end{aligned}$ | $\begin{aligned} & 50 \\ & 50 \\ & 50 \end{aligned}$ | $\begin{aligned} & 52 \\ & 52 \\ & 52 \\ & 52 \end{aligned}$ |
| $\stackrel{\omega_{\text {ABC }}}{ }$ | 0 | 0 |  | 0 | $317 \pm 6$ | 2.3 | $\leq 16$ | $\geq 9$ | $p d$ |  | $\pi^{+} \pi^{-}$ |  | 38 | 49 |

and DuMond, Nuovo Cimento 5, 541 (1957) and Fundamental Constants of Physics (Interscience Publishers, Inc., New York, 1957).
8. The neutron mean life is based on the value $11.7 \pm 0.3$ min for the half life, by A. N. Sosnovskij, P. E. Spivak, Y. A. Prokoviev, I. E. Kutikov, and Y. P. Dobrynin, Nucl. Phys. 10, 395 (1959).
9. $m\left(K^{+}\right)$is from the compilation of E. O. Okonov, Fortschr. Physik 8, 42 (1960), who also gives the $K^{-}$mass as 493.9 $\pm 0.4 \mathrm{MeV}$.
$\tau\left(K^{+}\right)$is a weighted average of the following results in $10^{-8} \mathrm{sec}: 1.224 \pm 0.013$, compilation of Barkas \& Rosenfeld (cf. reference 3), $1.231 \pm 0.011$, A. M. Boyarski, E. C. Loh, L. Q. Niemela, D. M. Ritson, R. Weinstein and S. Ozaki, Phys. Rev. 128, 2398 (1962).

The $K^{+}$branching ratios are from B. P. Roe, D. Sinclair, J. L. Brown, D. A. Glaser, J. A. Kadyk, and G. H. Trilling, Phys. Rev. Letters 7, 346 (1961); and G. Giacomelli, D. Monti, G. Quareni, A. Quareni-Vignudelli, W. Püschel, and J. Tietge, Phys. Letters 3, 346 (1963). Note that these branching ratios (except for $\tau$ ) disagree with the weighted averages obtained from emulsion experiments, as quoted by Crawford (cf. reference 6):

$$
\begin{aligned}
\mu 2 & =57.4 \pm 2.0 \\
\pi 2 & =25.6 \pm 1.5 \\
\mu 3+e 3+\tau & =11.0 \pm 1.0 \\
\tau & =5.7 \pm 0.2
\end{aligned}
$$

10. J. Button and B. C. Maglic, Phys. Rev. 127, 1297 (1962).
11. $m\left(K^{0}\right)$ is from A. H. Rosenfeld, F. T. Solmitz and R. D. Tripp, Phys. Rev. Letters 2, 110 (1959). $\mu\left(K^{0}\right)$ is from E. O. Okonov, J. Exptl. Theoret. Phys. (U.S.S.R.) 42, 1554 (1962).
$\tau\left(K_{1}^{0}\right)$ is a weighted average, based on the following recent results only (in $10^{-10} \mathrm{sec}$ ), all quoted by Crawford (cf. reference 6 ):
$0.94 \pm 0.05$, Crawford et al. (cf. below) ,
$0.90 \pm 0.05$, A. F. Garfinkel, Report Nevis 104 (1962). [Thesis, Columbia University, Physics Department (unpublished).]
$0.885 \pm 0.025$, R. L. Golden, G. Alexander, J. A. Anderson, F. S. and B. B. Crawford, L. J. Lloyd, G. W. Meisner, and L. Price ( to be published).

The $K_{1}^{0} \rightarrow \pi^{0} \pi^{0}$ branching ratio is a weighted average of the following results in percent:
$28.5 \pm 3.6$ by F. S. Crawford, M. Cresti, R. Douglass, M. Good, G. Kalbfleisch, M. L. Stevenson, and H. Ticho, Phys. Rev. Letters 2, 266 (1959),
$29.4 \pm 2.1$, M. Chretien, V. Fisher, H. R. Crouch, R. E. Lanou, J. T. Massimo, A. M. Shapiro, J. P. Averell, A. E. Brenner, D. R. Firth, L. G. Hyman, M. E. Law, R. H. Milburn, E. E. Ronat, K. Strauch, J. C. Street, J. J. Szymanski, L. Guerriero, I. A. Pless, L. Roesenson, and G. A. Saladin (to be published), quoted by Crawford (cf. reference 6),
$33.5 \pm 1.4$, J. L. Brown, J. A. Kadyk, G. H. Trilling, B. P. Roe, D. Sinclair, and J. C. Vander Welde, Phys. Rev. (to be published),
$26.0 \pm 2.4$, Anderson et al. (cf. reference 6),
$29 \pm 3$, weighted average of earlier results, reported at the 1960 Rochester Conference.
12. $\tau\left(K_{2}^{0}\right)$ is a weighted average of the following results in $10^{-8}$ sec
$6.8(+2.6 /-1.5)$, Crawford (cf. reference 6),
$8.1(+3.3 /-2.4)$, M. Bardon, K. Lande, L. M. Lederman, and W. Chinowsky, Ann. Phys. 5, 156 (1958).
$5.1(+2.4 /-1.3)$, S. E. Darmon, A. Rousset and W. Six, Phys. Letters 3, 57 (1962).

The branching ratios have been obtained from the following results:
$R_{1}=\frac{K_{2}^{0} \rightarrow 3 \pi^{0}}{K_{2}^{0} \rightarrow \text { all charged }}=0.38 \pm 0.07$, M. H. Anikina, M. S. Zhuravleva, D. M. Kotliarevsky, Z. S. Mandyavidze, A. M. Mestvirishvili, D. Neagu, E. O. Okonov, N. S. Petrov, A. M. Rosanova, V. A. Rusakov, G. G. Tachtamishev, and L. V. Chekhaidze, Proceedings of the 1962 Annual International Conference on High-Energy Physics at Geneva (CERN, Geneva, Switzerland, 1962), p. 452.
$R_{2}=\frac{K_{2}^{0} \rightarrow \pi^{+} \pi^{-} \pi^{0}}{K_{2}^{0} \rightarrow \text { all charged }}=0.127 \pm 0.020$, D. Luers, I. S. Mittra, W. J. Willis, and S. S. Yamamoto, Proceedings of the Aix-en-Provence International Conference on Elementary Particles in 1961 (CEN Saclay, Seine et Oise, France,1962), Vol. I, 241 (1961).
$R_{2}=0.134 \pm 0.018$, Anikina, et al.
$R_{3}=\frac{K_{2}^{0} \rightarrow \pi e \nu_{e}}{K_{2}^{0} \rightarrow \text { all charged }}=0.458 \pm 0.048$, Luers, et al.
$R_{3}=0.415 \pm 0.120$, A. Astier, L. Blaskovic, M. M. de Courreges, B. Equer, A. Lloret, P. Rivet, and J. Siaud, Proceedings of the Aix-en-Provence International Conference on Elementary Particles in 1961 (CEN Saclay, Seine et Oise, France, 1962), Vol. I, 227 (1961). A value on $R_{3}=\left(0.185_{-0.034}^{+0.038}\right)$ by Astier et al. has not been used.
13. $\tau\left(\pi^{0}\right)$ is from G. von Dardel, D. Dekkers, R. Mermod, J. D. van Putten, M. Vivargent, G. Weber, and K. Winter, Phys. Letters 4, 51 (1963).
The relative frequency for $\gamma e^{+} e^{-}$-decay is from J. Tietge and W. Püschel, Phys. Rev. 127, 1324 (1962).
14. $\mu\left(\mu^{+}\right)$is from G. Charpak, F. Farley, R. L. Garwin, T. Muller, J. C. Sens, and A. Zichichi, Phys. Letters 1, 16 (1962).
$m\left(\mu^{+}\right)$is a combined value of the latest measurements, as given by G. McD. Bingham, Nuovo Cimento 27, 1352 (1963). $\tau\left(\mu^{+}\right)$is a weighted average from R. A. Lundy, Phys. Rev. 125, 1686 (1962).
15. $\mu(e)$ and $\mu(p)$ are from A. A. Schupp, R. W. Pidd and H. R. Crane, Phys. Rev. 121, 1 (1961).
16. J. Bahcall and R. B. Curtis, Nuovo Cimento 21, 422 (1961).
17. L. M. Langer and R. J. D. Moffat, Phys. Rev. 88, 689 (1952).
18. $m\left(\pi^{ \pm}\right)$is from G. Shapiro and L. M. Lederman, Phys. Rev. 125, 1022 (1962),
$\tau\left(\pi^{ \pm}\right)$from A. W. Merrison, Advan. Phys. 11, 41, 1 (1962).
19. The parity of $\Sigma$ is + from measurements of R. D. Tripp, M. B. Watson and M. Ferro-Luzzi, Phys. Rev. Letters 8, 175 (1962), who conclude that $P(\Sigma p K)=-1$.
20. The $\Lambda$ parity is + from J. W. Cronin and O. E. Overseth, Proceedings of the 1962 Annual International Conference on High Energy Physics at Geneva (CERN, Geneva, Switzerland, 1962), p. 453, who give support for $P(K \Lambda)=-1$.
21. $\nu_{\mu}$ and $\nu_{e}$ are left-handed screw states, $\bar{\nu}_{\mu}$ and $\bar{\nu}_{e}$ righthanded. $\gamma$ has two spin states: right- and left-handed. We define a left-(right-)handed screw-state as a state with negative (positive) spin.
22. The $\pi^{0}$ mass has been obtained from the $\pi^{+}$mass (cf. reference 18) and the measurement by J. B. Czirr, University of California Radiation Laboratory Technical Report UCRL9951 (unpublished) and Bull. Am. Phys. Soc. 7, 265 (1962) of the mass difference $m\left(\pi^{-}\right)-m\left(\pi^{0}\right)=(4.6064 \pm 0.0030) \mathrm{MeV}$.
23. L. T. Kerth, Rev. Mod. Phys. 3, 389 (1961); O. Chamberlain, K. M. Crowe, D. Keefe, L. T. Kerth, A. Lemonick, Tin Maung, and T. F. Zipf, Phys. Rev. 125, 1696 (1962). E. F. Beall, W. Holley, D. Keefe, L. T. Kerth, J. J. Thresher, C. L. Wang, and W. A. Wenzel, Proceedings of the 1962 Annual International Conference on High Energy Physics at Geneva (CERN, Geneva, Switzerland, 1962), p. 368; E. V. Kuznetsov and Ya. Ya. Shalamov, J. Exptl. Theoret. Phys. (U.S.S.R.) 43, 1979 (1962).
24. V. V. Barmin, Y. S. Krestnikov, E. V. Kuznetsov, A. G. Meshkovsky, and V. A. Shebanov, J. Exptl. Theoret. Phys. (U.S.S.R.) 43, 1564 (1962); March et al. (cf. reference 57 ); Kuznetsov et al. (cf. reference 23).
25. R. H. March, A. R. Erwin, and W. D. Walker, Phys. Letters 3, 99 (1962);
W. Koch, J. D. Dowell, B. Leontic, A. Lundby, R. Meunier, J. P. Stroot, and M. Szeptycka, Phys. Letters 1, 53 (1962);
D. Colley, N. Gelfand, U. Nauenberg, J. Steinberger, S. Wolf, H. R. Brugger, P. R. Kramer, and R. J. Plano, Proceedings of the 1962 Annual International Conference on High Energy Physics at Geneva (CERN, Geneva, Switzerland, 1962), p. 315 and Phys. Rev. 128, 1930 (1962).
26. $m\left(\Xi_{1}^{*}\right)$ is a weighted average of the following results in MeV:
$1529 \pm 5$, G. M. Pjerrou, D. J. Prowse, P. Schlein, W. E. Slater, D. H. Stork, and H. K. Ticho, Phys. Rev. Letters 9, 114 (1962).
$1535 \pm 3$, L. Bertanza, V. Brisson, P. L. Connolly, E. L. Hart, I. S. Mittra, G. C. Moneti, R. R. Rau, N. P. Samios, I. O. Skillicorn, S. S. Yamamoto, M. Goldberg, L. Gray, J. Leitner, S. Lichtman, and J. Westgard, Phys. Rev. Letters 9, 180 (1962).
The spin and parity assignments are communicated privately by A. H. Rosenfeld.
27. M. Ferro-Luzzi, R. D. Tripp and M. B. Watson, Phys. Rev. Letters 8, 28 (1962). Further evidence has been produced by G. Alexander, G. R. Kalbfleisch, D. H. Miller, and D. A. Smith, Phys. Rev. Letters 8, 447 (1962); M. H. Alston, L. W. Alvarez, M. Ferro-Luzzi, A. H. Rosenfeld, H. K. Ticho, and S. G. Wojcicki, Proceedings of the 1962 Annual International Conference on High Energy Physics at Geneva (CERN, Geneva, Switzerland, 1962), p. 311; L. Bertanza, V. Brisson, P. L. Connolly, E. L. Hart, I. S. Mittra, G. C. Moneti, R. R. Rau, N. P. Samios, I. O. Skillicorn, S. S. Yamamoto, M. Goldberg, L. Gray, J. Leitner, S. Lichtman, and J. Westgard, ibid., p. 279; Bastien et al. (cf. reference 56); March et al. (cf. reference 25); Alexander et al. (cf. reference 56).
28. $m\left(Y_{0}^{*}\right)$ and $\Gamma\left(Y_{0}^{*}\right)$ are from $\AA$. Frisk and G. Ekspong, Phys. Letters 3, 27 (1962). Further evidence: Y. Eisenberg, G. Yekutieli, P. Abrahamson, and D. Kessler, Nuovo Cimento 21, 563 (1961); M. H. Alston, L. W. Alvarez, P. Eberhard, M. L. Good, W. Graziano, H. K. Ticho, and S. G. Wojcicki,

Phys. Rev. Letters 6, 698 (1961); P. Bastien, M. Ferro-Luzzi and A. H. Rosenfeld, Phys. Rev. Letters 6, 702 (1961); Y. International Conference on Elementary Particles in 1961 (CEN Saclay, Seine et Oise, France, 1962), p. 389; Wojcicki, University of California Radiation Laboratory Report UCRL-9835 (unpublished); Alston, et al. (cf. reference 27); Colley, et al. (cf. reference 25); Alexander, et al. (cf. reference 27 and 56); March et al. (cf. reference 57).
Note that all bubble chamber experiments are consistent with $\Gamma\left(Y_{0}^{*}\right)=50 \mathrm{MeV}$.
29. Review article by M. H. Alston and L. W. Ferro-Luzzi, Rev. Mod. Phys. 33, 416 (1961); M. Taher-Zadeh, D. J. Prowse, D. H. Stork, and H. K. Ticho, Bull. Am. Phys. Soc. 6, 510 (1961); M. M. Block, A. Engler, R. Gessaroli, J. Kopelman, M. Meer, A. Pevsner, P. Schlein, R. Strand, L. Grimellini, L. Lendinara, and L. Monari, Bull. Am. Phys. Soc. 7, 49 (1962); A. R. Erwin, R. H. March and W. D. Walker, (cf. reference 25); Nuovo Cimento 24, 237 (1962); Collaboration Saclay-Orsay-Bari-Bologne, Proceedings of the Aix-en-Provence International Conference on Elementary Particles in 1961 (CEN Saclay, Seine et Oise, France, 1962), p. 375; Dowell et al. (cf. reference 25); Eisenberg and Kessler (cf. reference 28); J. Auman, M. M. Block, R. Gessaroli, J. Kopelman, S. Ratti, L. Grimellini, T. Kikuchi, L. Lendinara, L. Monari, and E. Harth, Proceedings of the 1962 Annual International Conference on High Energy Physics at Geneva (CERN, Geneva, Switzerland, 1962), p. 330; J. Button-Shafer, M. Ferro-Luzzi, J. Murray, M. L. Stevenson, and F. T. Solmitz, ibid., pp. 303 and 307; C. T. Coffin, L. J. Curtis, D. I. Meyer, and K. M. Terwilliger, ibid., p. 327; Alexander et al. (cf. reference 27 and 56 ); Colley et al. (cf. reference 25); Bertanza et al. (cf. reference 27); Bastien et al. (cf. reference 56).

The spin and parity assignments are from J. B. Shafer, J. J. Murray, and D. O. Huwe, Phys. Rev. Letters 10, 179 (1963) and L. Bertanza, V. Brisson, P. L. Connolly, E. L. Hart, I. S. Mittra, G. C. Moneti, R. R. Rau, N. P. Samios, I. O. Skillicorn, S. S. Yamamoto, M. Goldberg, J. Leitner, S. Lichtman, and J. Westgard, Phys. Rev. Letters 10, 176(1963).
30. $Z_{3}^{*}$ is reported by Erwin et al. (cf. reference 29); $Z_{1}^{*}$ and $Z_{3}^{*}$ by Y. V. Kuznetsov, Ya. Ya. Shalamov, A. F. Grashin, and Y. P. Kuznetsov, Phys. Letters 1, 314 (1962), and $Z_{1}^{*}$ by A. I. Baz, V. G. Vaks and A. I. Larkin, Nucl. Phys. 38, 211 (1962), and L. Bertanza, P. L. Connolly, B. B. Culwick, F. R. Eisler, T. Morris, R. Palmer, A. Prodell, and A. Samios, Phys. Rev. Letters 8, 332 (1962). No evidence for $Z_{1}^{*}$ is found by Alexander et al. (cf. reference 56). $Z_{1}^{*}$ may turn out to be identical with $N_{15}^{*}$.
(A) in the reaction column denotes heavy nuclei. The momenta were $1.89 \mathrm{GeV} / c$ in the case of Erwin et al. (reference 29) and $2.8 \mathrm{GeV} / c$ in the case of Kuznetsov et al. The spin and isospin assignments seem to have been deduced theoretically.
31. P. Falk-Vairant and G. Valladas, Proceedings of the 1960 Annual International Conference on High Energy Physics at Rochester (University of Rochester, Rochester, 1960) p. 38, and Rev. Mod. Phys. 33, 362 (1961); B. J. Moyer, Rev. Mod. Phys. 33, 367 (1961); J. C. Brisson, P. Falk-Vairant, J. P. Merlo, P. Sonderegger, R. Turlay, and G. Valladas, Proceedings of the Aix-en-Provence International Conference on Elementary Particles in 1961, (CEN Saclay, Seine et Oise, France, 1962), Vol. I, p. 45 ; J. F. Detoeuf, ibid., II, p. 57; N. P. Samios, A. H. Bachman, R. M. Lea, T. E. Kalogeropoulos, and W. D. Shephard, Phys. Rev. Letters 9, 139 (1962); E. L. Hart, R. I. Louttit, D. Luers, T. W. Morris, W. J. Willis, and S. S. Yamamoto, Phys. Rev. 126, 747 (1962); Carmony et al. (cf. reference 37). According to J. A. Helland, T. J. Devlin, D. E. Hagge, M. J. Longo, B. J. Moyer and C. D. Wood, Bull. Am. Phys. Soc. 7, 468 (1962) and Phys. Rev. Letters 10, 27 (1963), there seems to be a superposition of $D_{5 / 2}$ and $F_{5 / 2}$ waves in the $N_{15}^{*}$ peak and a $J=7 / 2$ wave giving the largest contribution to the $N_{37}^{*}$, whereas no single state seems to give a dominant contribution to the $N_{13}^{*}$ peak. Cf. also T. J. Devlin, B. J. Moyer and V. Perez-Mendez, Phys. Rev. 125, 690 (1962). The $N_{15}^{*}$ and $N_{37}^{*}$ decay branching ratios are communicated privately by A. H. Rosenfeld.
32. B. T. Feld and W. M. Layson, Proceedings of the Annual 1962 International Conference on High Energy Physics at Geneva (CERN, Geneva, Switzerland, 1962) p. 147.
33. $m\left(f^{0}\right)$ is a weighted average of the following results in

MeV: $1260 \pm 35$, J. J. Veillet, J. Hennessy, H. Bingham, M. Block, D.' Drijard, A. Lagarrigue, P. Mittner, A. Rousset, G. Bellini, M. di Corato, E. Fiorini, and P. Negri, Phys. Rev. Letters 10, 29 (1963); $1250 \pm 25$, W. Selove, V. Hagopian, H. Brody, A. Baker, and E. Leboy, Phys. Rev. Letters 9, 272 (1962). Evidence against $f^{0}$ has been given by C. Alff, D. Berley, D. Colley, N. Gelfand, U. Nauenberg, D. Miller, J. Schultz, J. Steinberger, T. H. Tan, H. Brugger, P. Kramer, and R. Plano, Phys. Rev. Letters 9, 322 (1962).
34. $m\left(K_{1}^{*}\right)$ and $\Gamma\left(K_{1}^{*}\right)$ are "best values" from B. P. Gregory, Proceedings of the 1962 Annual International Conference on High-Energy Physics at Geneva (CERN, Geneva, Switzerland, 1962), p. 779. March et al. (cf. reference 25) give $\Gamma\left(K^{*}\right)$ $-(35 \pm 15) \mathrm{MeV}$, in agreement with this. The spin and parity assignments are from W. Chinowsky, G. Goldhaber, S. Goldhaber, W. Lee, and T. O'Halloran, Phys. Rev. Letters 9, 330 (1962).
The $K_{1}^{*-}$ branching ratios are weighted averages of the following results:
$\Gamma\left(\overline{K^{0}} \pi^{-}\right) / \Gamma\left(K^{-} \pi^{0}\right)=1.4 \pm 0.4, \mathrm{M}$. Alston, L. W. Alvarez, P. Eberhard, M. L. Good, W. Graziano, H. K. Ticho, and S. G. Wojcicki, Phys. Rev. Letters 6, 300 (1961).
$\Gamma\left(K^{-} \pi^{0}\right) / \Gamma\left(\bar{K}^{0} \pi^{-}\right)=0.5 \pm 0.2$, W. Graziano and S. G. Wojcicki, Phys. Rev. 128, 1868 (1962).
35. The $\pi^{-}$momentum for $K_{1}^{* 0}$ production is here tabulated as $1657 \mathrm{MeV} / c$, which corresponds to associated production with $\Lambda$. Associated production with $\Sigma^{0}$ gives $1826 \mathrm{MeV} / c$.
36. D. O. Caldwell, E. Bleuler, B. Elsner, L. W. Jones, and B. Zacharov, Phys. Letters 2, 253 (1962). Evidence against the $\varphi_{3}$ has been given by Alff et al. (cf. reference 33). A $T=2 \pi^{-} \pi^{-} \pi^{0}$ resonance at 870 MeV is indicated in the process $\pi^{-} n \rightarrow \pi^{-} \pi^{-} \pi^{0} p$ as reported by Shalamov et al. (cf. reference 51).
37. $m\left(\omega^{0}\right)$ is a weighted average of the following results in MeV:
$782 \pm 1$, Alff et al. (cf. reference 33).
$779.4 \pm 1.4$, R. Armenteros, R. Budde, L. Montanet, D. Morrison, S. Nilsson, A. Shapira, J. Vandermeulen, C. d'Andlau, A. Astier, C. Ghesquière, B. Gregory, D. Rahm, P. Rivet, and F. Solmitz, Proceedings of the 1962 Annual International Conference on High Energy Physics at Geneva (CERN, Geneva, Switzerland, 1962), p. 90.
The branching ratios are based on the upper limits given by G. R. Lynch, Proc. Phys. Soc. (London) 80, 46 (1962) and by C. Alff, D. Berley, D. Colley, N. Gelfand, U. Nauenberg, D. Miller, J. Schultz, J. Steinberger, T. H. Tan, H. Brugger, P. Kramer, and R. Plano, Phys. Rev. Letters 9, 325 (1962), and on the weighted average of the following results for $(\omega \rightarrow$ neutrals $) /\left(\omega \rightarrow \pi^{+} \pi^{-} \pi^{0}\right)$ :
$10 \pm 4$, Alff et al.
$21 \pm 7.5$, Armenteros et al.
$7 \pm 6$, M. Meer, R. Strand, R. Kraemer, L. Madansky, M. Nussbaum, A. Pevsner, C. Richardson, T. Toohig, M. Block, S. Orenstein, and T. Fields, Proceedings of the 1962 Annual International Conference on High-Energy Physics at Geneva (CERN, Geneva, Switzerland, 1962), p. 103.
$25 \pm 10$, J. Button-Shafer, M. Ferro-Luzzi, J. Murray, M. L. Stevenson, and F. T. Solmitz, ibid., p. 307.

By "neutrals" is primarily meant the $\pi^{0} \gamma$ channel. The $\pi^{+} \pi^{-}$branching ratio is communicated privately by A. H. Rosenfeld.
$\Gamma\left(\omega^{0}\right)$ is consistent with zero in all measurements. The lowest upper limit is 12 MeV , as given by M. L. Stevenson, L. W. Alvarez, B. C. Maglić, and A. H. Rosenfeld, Phys. Rev. 125, 687 (1962).
Further evidence has been given by B. C. Maglić, L. W. Alvarez, A. H. Rosenfeld, and M. L. Stevenson, Phys. Rev. Letters 7, 178 (1961); N.'H. Xuong and G. R. Lynch, Phys. Rev. Letters 7, 327 (1961); Nuovo Cimento 25, 923 (1962); Phys. Rev. 128, 1849 (1962); A. Pevsner, R. Kraemer, M. Nussbaum, C. Richardson, P. Schlein, R. Strand, T. Toohig, M. Block, A. Engler, R. Gessaroli, and C. Meltzer, Phys. Rev. Letters 7, 421 (1961); Hart et al. (cf. reference 31); W. D. Walker, E. West, A. R. Erwin, and R. H. March, Proceedings of the 1962 Annual International Conference on High-Energy Physics at Geneva (CERN, Geneva, Switzerland, 1962), p. 42; C. Richardson, R. Kraemer, M. Meer, M. Nussbaum, A.

Pevsner, R. Strand, T. Toohig, and M. Block, ibid., p. 96; T. Toohig, R. Kraemer, L. Madansky, M. Meer, M. Nussbaum, A. Pevsner, C. Richardson, R. Strand, and M. Block, ibid., p. 99; T. Ferbel, J. Sandweiss, H. D. Taft, M. Gailloud, T. W. Morris, R. M. Lea, and T. E. Kalogeropoulos, ibid., p. 76; D. D. Carmony, F. Grard, R. T. Van de Walle, and Nguyen-huu Xuong, ibid., p. 44; G. B. Chadwick, W. Davies, M. Derrick, C. Hawkins, P. B. Jones, J. H. Mulvey, D). Radojicic, C. A. Wilkinson, M. Cresti, A. Grigoletto, S. Limentani, A. Loria, L. Peruzzo, and R. Santangelo, ibid., p. 73.
38. $m\left(\rho^{-}\right)$and $\Gamma\left(\rho^{-}\right)$are weighted averages of the following results, tabulated as functions of the lab kinetic energy $T \pi$ :

| $\begin{aligned} & m\left(\rho^{ \pm}\right) \\ & (\mathrm{MeV}) \end{aligned}$ | $\begin{gathered} \Gamma\left(\rho^{ \pm}\right) \\ (\mathrm{MeV}) \end{gathered}$ | $\begin{gathered} T_{\pi} \\ (\mathrm{BeV}) \end{gathered}$ | References |
| :---: | :---: | :---: | :---: |
| $713 \pm 81$ | $31 \pm 143$ | 0.91 | Foelsche et al. (cf. reference 46) |
| $748 \pm 16$ |  | 1.1 | V. P. Kenney, W. D. Shephard and C. D. Gall, Phys. Rev. 126, 736 (1962) and Nuovo Cimento 23, 245 (1962) |
| $755 \pm 10$ | $61 \pm 24$ | 1.26 | Foelsche et al. (cf. reference 46) |
| $752 \pm 13$ |  | 1.45 | Saclay - Orsay - Bari -Bologna-collaboration, Aix $^{\text {a }}$, p. 257 and Nuovo Cimento (to be published) |
| $740( \pm 13)$ | $120( \pm 15)$ | 1.72-1.93 | Walker et al. (cf. reference 37 ) |
| $770 \pm 10$ | $130 \pm 10$ | $2.19-2.73$ | Alff et al. (cf. reference 33) |
| $780( \pm 25)$ |  | 7.0 | A. F. Grashin and Ya. Ya. Shalamov,CERN ${ }^{\text {b }}$, p. 58 |
| $755( \pm 15)$ | $110( \pm 15)$ | $\bar{p} p$ at rest | G. B. Chadwick, W. T. Davies, M. Derrick, C. Hawkins, J. H. Mulvey, D. Radojicic, C. A. Wilkinson, M. Cresti, S. Limentani, and R. Santangelo, CERN ${ }^{\text {b }}$, p. 69 and Phys. Rev. Letters 10, 62 (1963) |

[^1]The above mass values are all consistent with $757 \pm 5$, and no evidence for an energy dependence can be seen. In fact, the only such evidence has been presented by Foelsche et al., who give $m\left(\rho^{+}\right)=726 \pm 10$ and $\Gamma^{\prime}\left(\rho^{+}\right)=57 \pm 27$ at $T^{\prime} \pi=1.09$ BeV . This value has not been used by us.

Further evidence on $\rho^{-}$has been given by M. S. Aipnutdinov, S. Ya. Nikitin, Ya. M. Selector, A. F. Grashin and S. M. Zombkovskii, Proceedings of the 1962 Annual International Conference on High Energy Physics at Geneva (CERN, Geneva, Switzerland, 1962), p. 61 and Soviet Phys.-JETP 15, 979 (1962); A. F. Grashin and Ya. Ya. Shalamov, Soviet Phys.JETP 15, 770 and 787 (1962); W. I). Shephard and W. D). Walker, Phys. Rev. 126, 278 (1962); J. Derrick, quoted by Lynch, Proc. Phys. Soc. (London) 80, 46 (1962); E. West, J. Bishop, J. Boyd, A. R. Erwin, D. Lyon, R. H. March, P. H. Satterblom, and D. H. Walker, Bull. Am. Phys. Soc. 7, 281 (1962); Z. G. T. Guiragossian, W. M. Powell and H. S. White, Bull. Am. Phys. Soc. 7, 281 (1962); Grote, Klabuhn, Klugow, Krecker, Kundt, Lanius, and Meier, Nucl. Phys. 34, 648 and 659 (1962).
39. Further evidence on $\rho$ has been given by J. A. Anderson, V. X. Bang, P. G. Burke, D. D. Carmony, and N. Schmitz, Phys. Rev. Letters 6, 365 (1961); A. H. Rosenfeld, D. D.

Carmony and R. T. Van de Walle, Phys. Rev. Letters 8, 293 (1962); E. M. Friedlander, Phys. Rev. 127, 247 (1962).
40. $m\left(\rho^{0}\right)$ and $\Gamma\left(\rho^{0}\right)$ are weighted averages of the following results, tabulated as functions of the lab kinetic energy $T_{\pi}$ :

| $m\left(\rho^{0}\right)$ <br> $(\mathrm{MeV})$ | $\Gamma\left(\rho^{0}\right)$ <br> $(\mathrm{MeV})$ | $T_{\pi}$ <br> $(\mathrm{BeV})$ | References |
| :--- | :--- | :--- | :--- |
| $752 \pm 27$ |  | 1.45 | Saclay-Orsay-Bari-Bo- <br> logna Collaboration, <br> Nuovo Cimento (to be |
| $760( \pm 15)$ | 1.72 | published) <br> Walker et al. (cf. reference <br> 37) |  |
| $750 \pm 10$ | $100 \pm 10$ | $2.19-2.73$ | Alff et al. (cf. reference <br> $33)$ |
| $742( \pm 15)$ | 2.64 | Grashin et al. (cf. ref- <br> erence 38) |  |
| $747 \pm 17$ | $156 \pm 17$ | 4.55 | Samios et al. (cf. reference <br> 31) |
| $750( \pm 15)$ | 7.0 | Grote et al. (cf. reference <br> $38)$ |  |
| $755( \pm 15)$ | $110( \pm 15)$ | $\bar{p} p$ at restKenney et al. (cf. ref- <br> erence 38) |  |

There is no evidence for an energy dependence.
The branching ratios are from Meer et al. (cf. reference 37).
Further evidence on $\rho^{0}$ has been given by Maglic et al. (cf. reference 37); Guiragossian et al. (cf. reference 38); W. B. Johnson, L. B. Auerbach, T. Ypsilantis, C. E. Wiegand, J. Lach. and T. Elioff, Phys. Rev. Letters 9, 173 (1962); Shephard and Walker (cf. reference 38 ).
41. $m\left(\rho_{1}^{0}\right), m\left(\rho_{2}^{0}\right), \Gamma\left(\rho_{1}^{0}\right)$, and $\Gamma\left(\rho_{2}^{0}\right)$ are from J. Button, G. R. Kalbfleisch, G. R. Lynch, B. C. Maglić, A. H. Rosenfeld, and M. L. Stevenson, Phys. Rev. 126, 1858 (1962). Further evidence on $\rho_{1}^{\circ}$ and $\rho_{2}^{\circ}$ has been given by Walker et al. (cf. reference 37).
42. $m(\alpha)$ is given by E. Pickup, D. K. Robinson and E. O. Salant for $\alpha^{0}$ and $\alpha^{+}$in Phys. Rev. Letters 8, 329 (1962). An indication for $m\left(\alpha^{+}\right)=655 \mathrm{MeV}$ is given by B. Sechi Zorn in Phys. Rev. Letters 8, 282 (1962). Evidence against $\alpha$ has been given by Alff et al. (cf. reference 33). Indication for a $T=2\left(\pi^{-} \pi^{-} \pi^{0}\right)$ resonance at 630 MeV has been reported by Shalamov et al. (cf. reference 51).
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44. $m\left(\zeta^{-}\right)$is the weighted average given by Roos (cf. reference 43), where however the new values $(590 \pm 20) \mathrm{MeV}$ of R. Barloutaud, J. Heughebaert, A. Leveque, C. Louedec, J. Meyer, and D. Tycho, Nuovo Cimento 27, 238 (1963), and $(541 \pm 18) \mathrm{MeV}$ of the Saclay-Orsay-Bari-Bologna Collaboration, Nuovo Cimento (to be published), have been included. The last mentioned value is also used for $m\left(\zeta^{\circ}\right)$. The isospin assignment is from B. Sechi Zorn, Phys. Rev. Letters 8, 282 (1962).

Evidence against the existence of the $\zeta$ has been given by Alff et al. (cf. reference 33), Chadwick et al. (cf. reference 37); D. Stonehill, C. Baltay, H. Courant, W. Fickinger, E. C. Fowler, H. Kraybill, J. Sandweiss, J. Sanford and H. Taft, Phys. Rev. Letters 6, 624 (1961). Selove et al. (cf. reference 33), Caldwell et al. (cf. reference 36), Samios et al. (cf. reference 31), Johnson et al. (cf. reference 40); R. R. Crittenden, B. Musgrave and H. J. Martin, Bull. Am. Phys. Soc. 7, 468 (1962); J. Kirz, J. Schwartz, and R. D. Tripp, UCRL-10676 and Phys. Rev. (to be published).
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$548 \pm 1$, Alff et al. (cf. reference 33);
$548 \pm 1$, H. Foelsche, E. C. Fowler, H. L. Kraybill, J. R. Sanford, and D. Stonehill, Phys. Rev. Letters 9, 223 (1962), at $1090 \mathrm{MeV} / c$;
$551 \pm 2$, Foelsche et al., at $1260 \mathrm{MeV} / c$;
$546 \pm 4$, Pickup et al. (cf. reference 42);
P. L. Bastien, J. P. Berge, O. I. Dahl, M. Ferro-
$550 \pm 1.5$ Luzzi, D. H. Miller, J. J. Murray, A. H. Rosen-
$548 \pm 2\left\{\begin{array}{l}\text { feld, and M. B. Watson, Phys. Rev. Letters 8, }\end{array}\right.$
114 (1962), from different decay modes.
$\Gamma\left(\eta^{0}\right)$ is from Bastien et al. Alff et al. and Foelsche et al. give $\Gamma<10 \mathrm{MeV}$.
The isospin assignment is from D. D. Carmony, A. H. Rosenfeld and R. T. Van de Walle, Phys. Rev. Letters 8, 117 (1962), and the spin and parities from M. Chrétien, F. Bulos, H. R. Crouch, Jr., R. E. Lanou, J. T. Massimo, A. M. Shapiro, J. A. Averell, C. A. Bordner, Jr., A. E. Brenner, D. R. Firth, M. E. Law, E. E. Ronat, K. Strauch, J. C. Street, J. J. Szymanski, A. Weinberg, B. Nelson, I. A. Pless, L. Rosenson, G. A. Salandin, R. K. Yamamoto, L. Guerriero, and F. Waldner, Phys. Rev. Letters 9, 127 (1962).
The branching ratios have been computed from the following measurements:
$R_{1}=\frac{\pi^{+} \pi^{-} \gamma}{\pi^{+} \pi^{-} \pi^{0}}=0.26 \pm 0.08$, E. C. Fowler, F. S. Crawford, L. J. Lloyd, R. A. Grossman, L. Price, Phys. Rev. Letters 10, 110 (1963).
$R_{2}=\frac{\text { neutrals }}{\pi^{+} \pi^{-} \pi^{0}}=2.5 \pm 0.5$, Alff et al.
$R_{2}=2.5 \pm 1.0$, Pickup et al.
$1 / R_{2}=0.31 \pm 0.11$, Bastien et al.
$R_{2} /\left(1+R_{1}\right)=2.7 \pm 0.8$, Button-Shafer et al. (cf. reference 29).
$R_{2} /\left(1+R_{1}\right)=3.1 \pm 1.2$, Meer et al. (cf. reference 37 ).
Chrétien et al. also give the branching ratio $R_{3}=\frac{3 \pi^{0}}{2 \gamma}$ $\leq 1.1 \pm 0.3$.
Further evidence on $\eta$ has been given by Pevsner et al. (cf. reference 37); Foster et al. (cf. reference 45); C. Mencuccini, R. Querzoli, G. Salvini, and V. G. Silvestrini, Proceedings of the 1962 Annual International Conference on HighEnergy Physics at Geneva (CERN, Geneva, Switzerland, 1962), p. 33; H. Foelsche, E. C. Fowler, H. L. Kraybill, J. R. Sanford, and D. Stonehill, the same conference, p. 36; W. D. Walker, J. Boyd, A. R. Erwin, H. R. Fechter, D. Lyon, R. H. March, P. H. Satterblom, and E. West, Bull. Am. Phys. Soc. 7, 281 (1962).
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$322 \pm 8$, B. Richter, Phys. Rev. Letters 9, 217 (1962).
Further evidence has been given by Button et al. (cf. reference 40 ). No evidence has been found by L. Lapidus, Proceedings of the 1962 Annual International Conference on HighEnergy Physics at Geneva (CERN, Geneva, Switzerland, 1962), p. 115 nor by Kirz et al. (cf. reference 44) nor by Blokhintseva et al. (cf. reference 47).
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[^0]:    * "Not known", here and henceforth is short for "not known to the compiler."

[^1]:    a Proceedings of the Aix-en-Provence International Conference on Elementary Particles in 1961 (CEN Saclay, Seine et Oise, France, 1962), Vol. I, (1961).
    b Proceedings of the Annual 1962 International Conference on High-Energy Physics at Geneva (CERN, Geneva, Switzerland, 1962).

