


 $I(J^P) = \frac{1}{2}(\frac{1}{2}^+)$ Status: ****

The parity has not actually been measured, but + is of course expected.

Ξ^0 MASS

The fit uses the Ξ^0 , Ξ^- , and Ξ^+ masses and the $\Xi^- - \Xi^0$ mass difference. It assumes that the Ξ^- and Ξ^+ masses are the same.

VALUE (MeV)	EVTS	DOCUMENT ID	TECN	COMMENT
1314.86±0.20 OUR FIT				
1314.82±0.06±0.20	3120	FANTI	00	NA48 p Be, 450 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •				
1315.2 ± 0.92	49	WILQUET	72	HLBC
1313.4 ± 1.8	1	PALMER	68	HBC

$m_{\Xi^-} - m_{\Xi^0}$

The fit uses the Ξ^0 , Ξ^- , and Ξ^+ masses and the $\Xi^- - \Xi^0$ mass difference. It assumes that the Ξ^- and Ξ^+ masses are the same.

VALUE (MeV)	EVTS	DOCUMENT ID	TECN	COMMENT
6.85±0.21 OUR FIT				
6.3 ± 0.7 OUR AVERAGE				
6.9 ± 2.2	29	LONDON	66	HBC
6.1 ± 0.9	88	PJERROU	65B	HBC
6.8 ± 1.6	23	JAUNEAU	63	FBC
• • • We do not use the following data for averages, fits, limits, etc. • • •				
6.1 ± 1.6	45	CARMONY	64B	HBC See PJERROU 65B

Ξ^0 MEAN LIFE

VALUE (10^{-10} s)	EVTS	DOCUMENT ID	TECN	COMMENT
2.90±0.09 OUR AVERAGE				
2.83±0.16	6300	¹ ZECH	77	SPEC Neutral hyperon beam
2.88 ^{+0.21} _{-0.19}	652	BALTAY	74	HBC 1.75 GeV/c $K^- p$
2.90 ^{+0.32} _{-0.27}	157	² MAYEUR	72	HLBC 2.1 GeV/c K^-
3.07 ^{+0.22} _{-0.20}	340	DAUBER	69	HBC
3.0 ± 0.5	80	PJERROU	65B	HBC
2.5 ^{+0.4} _{-0.3}	101	HUBBARD	64	HBC
3.9 ^{+1.4} _{-0.8}	24	JAUNEAU	63	FBC
• • • We do not use the following data for averages, fits, limits, etc. • • •				
3.5 ^{+1.0} _{-0.8}	45	CARMONY	64B	HBC See PJERROU 65B

¹The ZECH 77 result is $\tau_{\Xi^0} = [2.77 - (\tau_\Lambda - 2.69)] \times 10^{-10}$ s, in which we use $\tau_\Lambda = 2.63 \times 10^{-10}$ s.

²The MAYEUR 72 value is modified by the erratum.

Ξ^0 MAGNETIC MOMENT

See the “Quark Model” review.

VALUE (μ_N)	EVTS	DOCUMENT ID	TECN
-1.250 ± 0.014 OUR AVERAGE			
-1.253 ± 0.014	270k	COX	81 SPEC
-1.20 ± 0.06	42k	BUNCE	79 SPEC

Ξ^0 DECAY MODES

Mode	Fraction (Γ_i/Γ)	Confidence level
$\Gamma_1 \Lambda \pi^0$	$(99.524 \pm 0.012) \%$	
$\Gamma_2 \Lambda \gamma$	$(1.17 \pm 0.07) \times 10^{-3}$	
$\Gamma_3 \Lambda e^+ e^-$	$(7.6 \pm 0.6) \times 10^{-6}$	
$\Gamma_4 \Sigma^0 \gamma$	$(3.33 \pm 0.10) \times 10^{-3}$	
$\Gamma_5 \Sigma^+ e^- \bar{\nu}_e$	$(2.52 \pm 0.08) \times 10^{-4}$	
$\Gamma_6 \Sigma^+ \mu^- \bar{\nu}_\mu$	$(2.33 \pm 0.35) \times 10^{-6}$	

$\Delta S = \Delta Q$ (SQ) violating modes or
 $\Delta S = 2$ forbidden (S2) modes

$\Gamma_7 \Sigma^- e^+ \nu_e$	SQ	< 1.6	$\times 10^{-4}$	90%
$\Gamma_8 \Sigma^- \mu^+ \nu_\mu$	SQ	< 9	$\times 10^{-4}$	90%
$\Gamma_9 p \pi^-$	S2	< 8	$\times 10^{-6}$	90%
$\Gamma_{10} p e^- \bar{\nu}_e$	S2	< 1.3	$\times 10^{-3}$	
$\Gamma_{11} p \mu^- \bar{\nu}_\mu$	S2	< 1.3	$\times 10^{-3}$	
$\Gamma_{12} K^- e^+$				
$\Gamma_{13} K^+ e^-$				

CONSTRAINED FIT INFORMATION

An overall fit to 5 branching ratios uses 11 measurements and one constraint to determine 5 parameters. The overall fit has a $\chi^2 = 7.5$ for 7 degrees of freedom.

The following *off-diagonal* array elements are the correlation coefficients $\langle \delta x_i \delta x_j \rangle / (\delta x_i \cdot \delta x_j)$, in percent, from the fit to the branching fractions, $x_i \equiv \Gamma_i / \Gamma_{\text{total}}$. The fit constrains the x_i whose labels appear in this array to sum to one.

x_2	-57			
x_4	-82	0		
x_5	-7	0	0	
x_6	0	0	0	1

$x_1 \quad x_2 \quad x_4 \quad x_5$

Ξ^0 BRANCHING RATIOS

$\Gamma(\Lambda\gamma)/\Gamma(\Lambda\pi^0)$

<u>VALUE (units 10^{-3})</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
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1.17 ± 0.07 OUR FIT

1.17 ± 0.07 OUR AVERAGE

$1.17 \pm 0.05 \pm 0.06$	672	¹ LAI	04A	NA48	p Be, 450 GeV
$1.91 \pm 0.34 \pm 0.19$	31	² FANTI	00	NA48	p Be, 450 GeV
$1.06 \pm 0.12 \pm 0.11$	116	JAMES	90	SPEC	FNAL hyperons

¹ LAI 04A used our 2002 value of 99.5% for the $\Xi^0 \rightarrow \Lambda\pi^0$ branching fraction to get $\Gamma(\Xi^0 \rightarrow \Lambda\gamma)/\Gamma_{\text{total}} = (1.16 \pm 0.05 \pm 0.06) \times 10^{-3}$. We adjust slightly to go back to what was directly measured.

² FANTI 00 used our 1998 value of 99.5% for the $\Xi^0 \rightarrow \Lambda\pi^0$ branching fraction to get $\Gamma(\Xi^0 \rightarrow \Lambda\gamma)/\Gamma_{\text{total}} = (1.90 \pm 0.34 \pm 0.19) \times 10^{-3}$. We adjust slightly to go back to what was directly measured.

$\Gamma(\Lambda e^+ e^-)/\Gamma_{\text{total}}$

<u>VALUE (units 10^{-6})</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
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$7.6 \pm 0.4 \pm 0.5$	397 ± 21	¹ BATLEY	07C	NA48	p Be, 400 GeV
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¹ This BATLEY 07C result is consistent with internal bremsstrahlung.

$\Gamma(\Sigma^0\gamma)/\Gamma(\Lambda\pi^0)$

<u>VALUE (units 10^{-3})</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
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3.35 ± 0.10 OUR FIT

3.35 ± 0.10 OUR AVERAGE

$3.34 \pm 0.05 \pm 0.09$	4045	ALAVI-HARATI01C	KTEV	p nucleus, 800 GeV
$3.16 \pm 0.76 \pm 0.32$	17	¹ FANTI	00	NA48
$3.56 \pm 0.42 \pm 0.10$	85	TEIGE	89	SPEC FNAL hyperons

¹ FANTI 00 used our 1998 value of 99.5% for the $\Xi^0 \rightarrow \Lambda\pi^0$ branching fraction to get $\Gamma(\Xi^0 \rightarrow \Sigma^0\gamma)/\Gamma_{\text{total}} = (3.14 \pm 0.76 \pm 0.32) \times 10^{-3}$. We adjust slightly to go back to what was directly measured.

$\Gamma(\Sigma^+\bar{e}^-)/\Gamma_{\text{total}}$

<u>VALUE (units 10^{-4})</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
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2.52 ± 0.08 OUR FIT

2.53 ± 0.08 OUR AVERAGE

$2.51 \pm 0.03 \pm 0.09$	6101	BATLEY	07	NA48	p Be, 400 GeV
$2.55 \pm 0.14 \pm 0.10$	419	¹ BATLEY	07	NA48	p Be, 400 GeV
$2.71 \pm 0.22 \pm 0.31$	176	AFFOLDER	99	KTEV	p nucleus, 800 GeV

¹ This BATLEY 07 result is for $\Xi^0 \rightarrow \Sigma^- e^+ \nu_e$ events.

Γ_2/Γ_1

Γ_3/Γ

Γ_4/Γ_1

Γ_5/Γ

$\Gamma(\Sigma^+ \mu^- \bar{\nu}_\mu)/\Gamma_{\text{total}}$ Γ_6/Γ

<u>VALUE (units 10^{-6})</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
2.3 \pm 0.4 OUR FIT				
2.17 \pm 0.32 \pm 0.17	66	¹ BATLEY	13	NA48 p Be, 400 GeV

¹ BATLEY 13 used $\Xi^0 \rightarrow \Sigma^+ e^- \bar{\nu}_e$ decay as a normalization mode and its branching fraction value of $(2.51 \pm 0.03 \pm 0.09) \times 10^{-4}$ from BATLEY 07.

 $\Gamma(\Sigma^+ \mu^- \bar{\nu}_\mu)/\Gamma(\Sigma^+ e^- \bar{\nu}_e)$ Γ_6/Γ_5

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
0.0092 \pm 0.0015 OUR FIT				
0.018 \pm 0.007 \pm 0.002	9	ABOUZAID	05	KTEV p nucleus 800 GeV

 $\Gamma(\Sigma^- e^+ \nu_e)/\Gamma_{\text{total}}$ Γ_7/Γ

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<1.6 $\times 10^{-4}$	90	ABLIKIM	23B	BES3 $J/\psi \rightarrow \Xi^0 \Xi^0$

 $\Gamma(\Sigma^- e^+ \nu_e)/\Gamma(\Lambda \pi^0)$ Γ_7/Γ_1

Test of $\Delta S = \Delta Q$ rule.

<u>VALUE (units 10^{-3})</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
• • • We do not use the following data for averages, fits, limits, etc. • • •				
<0.9	90	YEH	74	HBC Effective denom.=2500
<1.5		DAUBER	69	HBC
<6		HUBBARD	66	HBC

 $\Gamma(\Sigma^- \mu^+ \nu_\mu)/\Gamma(\Lambda \pi^0)$ Γ_8/Γ_1

Test of $\Delta S = \Delta Q$ rule.

<u>VALUE (units 10^{-3})</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<0.9	90	YEH	74	HBC Effective denom.=2500
• • • We do not use the following data for averages, fits, limits, etc. • • •				
<1.5		DAUBER	69	HBC
<6		HUBBARD	66	HBC

 $\Gamma(p\pi^-)/\Gamma(\Lambda \pi^0)$ Γ_9/Γ_1

$\Delta S=2$. Forbidden in first-order weak interaction.

<u>VALUE (units 10^{-6})</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
< 8.2	90	WHITE	05	HYCP p Cu, 800 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •				
< 36	90	GEWENIGER	75	SPEC
<1800	90	YEH	74	HBC Effective denom.=1300
< 900		DAUBER	69	HBC
<5000		HUBBARD	66	HBC

 $\Gamma(p e^- \bar{\nu}_e)/\Gamma(\Lambda \pi^0)$ Γ_{10}/Γ_1

$\Delta S=2$. Forbidden in first-order weak interaction.

<u>VALUE (units 10^{-3})</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<1.3		DAUBER	69	HBC
• • • We do not use the following data for averages, fits, limits, etc. • • •				
<3.4	90	YEH	74	HBC Effective denom.=670
<6		HUBBARD	66	HBC

$\Gamma(p\mu^-\bar{\nu}_\mu)/\Gamma(\Lambda\pi^0)$ $\Delta S=2$. Forbidden in first-order weak interaction.

<u>VALUE</u> (units 10^{-3})	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<1.3		DAUBER	69	HBC
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$				
<3.5	90	YEH	74	HBC Effective denom.=664
<6		HUBBARD	66	HBC

 Γ_{11}/Γ_1 $\Gamma(K^- e^+)/\Gamma_{\text{total}}$ Γ_{12}/Γ

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<3.6 \times 10^{-6}$	90	¹ ABLIKIM	23AP BES3	$J/\psi \rightarrow \Xi^0 \Xi^0$

¹ This decay mode violates baryon and lepton number conservation with $\Delta(B-L) = 0$. $\Gamma(K^+ e^-)/\Gamma_{\text{total}}$ Γ_{13}/Γ

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<1.9 \times 10^{-6}$	90	¹ ABLIKIM	23AP BES3	$J/\psi \rightarrow \Xi^0 \Xi^0$

¹ This decay mode violates baryon and lepton number conservation with $|\Delta(B-L)|=2$. Ξ^0 DECAY PARAMETERS

See the “Note on Baryon Decay Parameters” in the neutron Listings.

 $\alpha(\Xi^0) \alpha_-(\Lambda)$ This is a product of the $\Xi^0 \rightarrow \Lambda\pi^0$ and $\Lambda \rightarrow p\pi^-$ asymmetries.

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
-0.261 ± 0.006 OUR AVERAGE				
-0.276 $\pm 0.001 \pm 0.035$	4M	BATLEY	10B	NA48 p Be, 400 GeV
-0.260 $\pm 0.004 \pm 0.005$	300k	HANDLER	82	SPEC FNAL hyperons
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$				
-0.317 ± 0.027	6075	BUNCE	78	SPEC FNAL hyperons
-0.35 ± 0.06	505	BALTAY	74	HBC $K^- p$ 1.75 GeV/c
-0.28 ± 0.06	739	DAUBER	69	HBC $K^- p$ 1.7–2.6 GeV/c

 α FOR $\Xi^0 \rightarrow \Lambda\pi^0$ The PDG 22 average $\alpha(\Xi^0)\alpha_-(\Lambda) = -0.261 \pm 0.006$ divided by the PDG 22 average $\alpha_-(\Lambda) = 0.748 \pm 0.007$ gives $\alpha(\Xi^0) = -0.349 \pm 0.009$.

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
-0.349 ± 0.009 OUR EVALUATION				
$-0.3750 \pm 0.0034 \pm 0.0016$	327k	ABLIKIM	23AU BES3	$J/\psi \rightarrow \Xi^0 \Xi^0$
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$				
-0.358 $\pm 0.042 \pm 0.013$	1.9k	ABLIKIM	23AD BES3	$\psi(2S) \rightarrow \Xi^0 \Xi^0$

 α FOR $\Xi^0 \rightarrow \bar{\Lambda}\pi^0$

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$0.3790 \pm 0.0034 \pm 0.0021$	327k	ABLIKIM	23AU BES3	$J/\psi \rightarrow \Xi^0 \Xi^0$
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$				
0.363 $\pm 0.042 \pm 0.013$	1.9k	ABLIKIM	23AD BES3	$\psi(2S) \rightarrow \Xi^0 \Xi^0$

ϕ ANGLE FOR $\Xi^0 \rightarrow \Lambda\pi^0$ ($(\tan\phi = \beta/\gamma)$)

VALUE (°)	EVTS	DOCUMENT ID	TECN	COMMENT
0.292 ± 0.550 ± 0.103	327k	¹ ABLIKIM	23AU BES3	$J/\psi \rightarrow \Xi^0 \Xi^0$
• • • We do not use the following data for averages, fits, limits, etc. • • •				
1.55 ± 6.70 ± 0.63	1.9k	ABLIKIM	23AD BES3	$\psi(2S) \rightarrow \Xi^0 \Xi^0$
16 ± 17	652	BALTAY	74	HBC 1.75 GeV/c $K^- p$
38 ± 19	739	² DAUBER	69	HBC
- 8 ± 30	146	³ BERGE	66	HBC

¹ Converted from radians to degrees.

² DAUBER 69 uses $\alpha_\Lambda = 0.647 \pm 0.020$.

³ The errors have been multiplied by 1.2 due to approximations used for the Ξ polarization; see DAUBER 69 for a discussion.

ϕ ANGLE FOR $\Xi^0 \rightarrow \bar{\Lambda}\pi^0$ with $\tan\phi = \beta/\gamma$

VALUE (degrees)	EVTS	DOCUMENT ID	TECN	COMMENT
- 0.304 ± 0.556 ± 0.109	327k	ABLIKIM	23AU BES3	$J/\psi \rightarrow \Xi^0 \Xi^0$
• • • We do not use the following data for averages, fits, limits, etc. • • •				
- 10.60 ± 6.65 ± 0.97	1.9k	ABLIKIM	23AD BES3	$\psi(2S) \rightarrow \Xi^0 \Xi^0$

$\Delta\phi_{CP}(\Xi^0) = (\phi_{\Xi^0} + \phi_{\Xi^0})/2$

VALUE (degrees)	EVTS	DOCUMENT ID	TECN	COMMENT
- 0.006 ± 0.395 ± 0.052	327k	¹ ABLIKIM	23AU BES3	$J/\psi \rightarrow \Xi^0 \Xi^0$
• • • We do not use the following data for averages, fits, limits, etc. • • •				
- 4.53 ± 4.70 ± 0.57	1.9k	¹ ABLIKIM	23AD BES3	$\psi(2S) \rightarrow \Xi^0 \Xi^0$

¹ Converted from radians to degrees.

A_{CP} FOR $\Xi^0 \rightarrow \Lambda\pi^0, \Xi^0 \rightarrow \bar{\Lambda}\pi^0$

$$A_{CP} = (\alpha + \bar{\alpha}) / (\alpha - \bar{\alpha})$$

VALUE (units 10^{-3})	EVTS	DOCUMENT ID	TECN	COMMENT
- 5.4 ± 6.5 ± 3.1	327k	¹ ABLIKIM	23AU BES3	$J/\psi \rightarrow \Xi^0 \Xi^0$
• • • We do not use the following data for averages, fits, limits, etc. • • •				
- 7 ± 82 ± 25	1.9k	¹ ABLIKIM	23AD BES3	$\psi(2S) \rightarrow \Xi^0 \Xi^0$

¹ Probing CP symmetry and weak phases with entangled double-strange baryons.

RADIATIVE HYPERON DECAYS

Revised July 2011 by J.D. Jackson (LBNL).

The weak radiative decays of spin-1/2 hyperons, $B_i \rightarrow B_f \gamma$, yield information about matrix elements (form factors) similar to that gained from weak hadronic decays. For a polarized spin-1/2 hyperon decaying radiatively via a $\Delta Q = 0$, $\Delta S = 1$ transition, the angular distribution of the direction \hat{p} of the final spin-1/2 baryon in the hyperon rest frame is

$$\frac{dN}{d\Omega} = \frac{N}{4\pi} (1 + \alpha_\gamma \mathbf{P}_i \cdot \hat{\mathbf{p}}) . \quad (1)$$

Here \mathbf{P}_i is the polarization of the decaying hyperon, and α_γ is the asymmetry parameter. In terms of the form factors $F_1(q^2)$, $F_2(q^2)$, and $G(q^2)$ of the effective hadronic weak electromagnetic vertex,

$$F_1(q^2)\gamma_\lambda + iF_2(q^2)\sigma_{\lambda\mu}q^\mu + G(q^2)\gamma_\lambda\gamma_5 ,$$

α_γ is

$$\alpha_\gamma = \frac{2 \operatorname{Re}[G(0)F_M^*(0)]}{|G(0)|^2 + |F_M(0)|^2} , \quad (2)$$

where $F_M = (m_i - m_f)[F_2 - F_1/(m_i + m_f)]$. If the decaying hyperon is unpolarized, the decay baryon has a longitudinal polarization given by $P_f = -\alpha_\gamma$ [1].

The angular distribution for the weak hadronic decay, $B_i \rightarrow B_f\pi$, has the same form as Eq. (1), but of course with a different asymmetry parameter, α_π . Now, however, if the decaying hyperon is unpolarized, the decay baryon has a longitudinal polarization given by $P_f = +\alpha_\pi$ [2,3]. The difference of sign is because the spins of the pion and photon are different.

$\Xi^0 \rightarrow \Lambda\gamma$ decay—The radiative decay $\Xi^0 \rightarrow \Lambda\gamma$ of an unpolarized Ξ^0 uses the hadronic decay $\Lambda \rightarrow p\pi^-$ as the analyzer. As noted above, the longitudinal polarization of the Λ will be $P_\Lambda = -\alpha_{\Xi\Lambda\gamma}$. Let α_- be the $\Lambda \rightarrow p\pi^-$ asymmetry parameter and $\theta_{\Lambda p}$ be the angle, as seen in the Λ rest frame, between the Λ line of flight and the proton momentum. Then the hadronic version of Eq. (1) applied to the $\Lambda \rightarrow p\pi^-$ decay gives

$$\frac{dN}{d\cos\theta_{\Lambda p}} = \frac{N}{2} (1 - \alpha_{\Xi\Lambda\gamma} \alpha_- \cos\theta_{\Lambda p}) \quad (3)$$

for the angular distribution of the proton in the Λ frame. Our current value, from the CERN NA48/1 experiment [4], is $\alpha_{\Xi\Lambda\gamma} = -0.704 \pm 0.019 \pm 0.064$.

$\Xi^0 \rightarrow \Sigma^0\gamma$ decay—The asymmetry parameter here, $\alpha_{\Xi\Sigma\gamma}$, is measured by following the decay chain $\Xi^0 \rightarrow \Sigma^0\gamma$, $\Sigma^0 \rightarrow \Lambda\gamma$, $\Lambda \rightarrow p\pi^-$. Again, for an unpolarized Ξ^0 , the longitudinal polarization of the Σ^0 will be $P_\Sigma = -\alpha_{\Xi\Sigma\gamma}$. In the $\Sigma^0 \rightarrow \Lambda\gamma$ decay, a parity-conserving magnetic-dipole transition, the polarization of the Σ^0 is transferred to the Λ , as may be seen as follows. Let $\theta_{\Sigma\Lambda}$ be the angle seen in the Σ^0 rest frame between the Σ^0 line of flight and the Λ momentum. For Σ^0 helicity $+1/2$, the probability amplitudes for positive and negative spin states of the Σ^0 *along the Λ momentum* are $\cos(\theta_{\Sigma\Lambda}/2)$ and $\sin(\theta_{\Sigma\Lambda}/2)$. Then the amplitude for a negative helicity photon and a negative helicity Λ is $\cos(\theta_{\Sigma\Lambda}/2)$, while the amplitude for positive helicities for the photon and Λ is $\sin(\theta_{\Sigma\Lambda}/2)$. For Σ^0 helicity $-1/2$, the amplitudes are interchanged. If the Σ^0 has longitudinal polarization P_Σ , the probabilities for Λ helicities $\pm 1/2$ are therefore

$$p(\pm 1/2) = \frac{1}{2}(1 \mp P_\Sigma) \cos^2(\theta_{\Sigma\Lambda}/2) + \frac{1}{2}(1 \pm P_\Sigma) \sin^2(\theta_{\Sigma\Lambda}/2), \quad (4)$$

and the longitudinal polarization of the Λ is

$$P_\Lambda = -P_\Sigma \cos \theta_{\Sigma\Lambda} = +\alpha_{\Xi\Sigma\gamma} \cos \theta_{\Sigma\Lambda}. \quad (5)$$

Using Eq. (1) for the $\Lambda \rightarrow p\pi^-$ decay again, we get for the joint angular distribution of the $\Sigma^0 \rightarrow \Lambda\gamma$, $\Lambda \rightarrow p\pi^-$ chain,

$$\frac{d^2N}{d \cos \theta_{\Sigma\Lambda} d \cos \theta_{\Lambda p}} = \frac{N}{4} (1 + \alpha_{\Xi\Sigma\gamma} \cos \theta_{\Sigma\Lambda} \alpha_- \cos \theta_{\Lambda p}). \quad (6)$$

Our current average for $\alpha_{\Xi\Sigma\gamma}$ is -0.69 ± 0.06 [4,5].

References

1. R.E. Behrends, Phys. Rev. **111**, 1691 (1958); see Eq. (7) or (8).
2. In ancient times, the signs of the asymmetry term in the angular distributions of radiative and hadronic decays of polarized hyperons were sometimes opposite. For roughly 50 years, however, the overwhelming convention has been to make them the same. The aim, not always achieved, is to remove ambiguities.
3. For the definition of α_π , see the note on “Baryon Decay Parameters” in the Neutron Listings.
4. J.R. Batley *et al.*, Phys. Lett. **B693**, 241 (2010).
5. A. Alavi-Harati *et al.*, Phys. Rev. Lett. **86**, 3239 (2001).

α FOR $\Xi^0 \rightarrow \Lambda\gamma$

See the note above on “Radiative Hyperon Decays.”

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
-0.704±0.019±0.064	52k	¹ BATLEY	10B NA48	p Be, 400 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •				
-0.78 ± 0.18 ± 0.06	672	LAI	04A NA48	See BATLEY 10B
-0.43 ± 0.44	87	² JAMES	90 SPEC	FNAL hyperons

¹ BATLEY 10B also measured the $\Xi^0 \rightarrow \bar{\Lambda}\gamma$ asymmetry to be -0.798 ± 0.064 (no systematic error given) with 4769 events.

² The sign has been changed; see the erratum, JAMES 02.

α FOR $\Xi^0 \rightarrow \Lambda e^+ e^-$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
-0.8±0.2	397 ± 21	¹ BATLEY	07C NA48	p Be, 400 GeV

¹ This BATLEY 07C result is consistent with the asymmetry α for $\Xi^0 \rightarrow \Lambda\gamma$, as expected if the mechanism is internal bremsstrahlung.

α FOR $\Xi^0 \rightarrow \Sigma^0\gamma$

See the note above on “Radiative Hyperon Decays.”

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
-0.69 ±0.06 OUR AVERAGE				
• • • We do not use the following data for averages, fits, limits, etc. • • •				
-0.729±0.030±0.076	15k	¹ BATLEY	10B NA48	p Be, 400 GeV
-0.63 ± 0.08 ± 0.05	4045	ALAVI-HARATI01C	KTEV	p nucleus, 800 GeV

+0.20 ± 0.32 ± 0.05 85 ² TEIGE 89 SPEC FNAL hyperons

¹ BATLEY 10B also measured the $\Xi^0 \rightarrow \bar{\Sigma}^0\gamma$ asymmetry to be -0.786 ± 0.104 (no systematic error given) with 1404 events.

² This result has been withdrawn, due to an error. See the erratum, TEIGE 02.

$g_1(0)/f_1(0)$ FOR $\Xi^0 \rightarrow \Sigma^+ e^- \bar{\nu}_e$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
1.22±0.05 OUR AVERAGE				
1.21±0.05		BATLEY	13	NA48 p Be, 400 GeV
1.32 ^{+0.21} _{-0.17} ±0.05	487	¹ ALAVI-HARATI01I	KTEV	p nucleus, 800 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •				
1.20±0.04±0.03	6520	² BATLEY	07	NA48 See BATLEY 13

¹ ALAVI-HARATI 01I assumes here that the second-class current is zero and that the weak-magnetism term takes its exact SU(3) value.

² This BATLEY 07 result uses our 2006 value of V_{us} from semileptonic kaon decays as input.

 $g_2(0)/f_1(0)$ FOR $\Xi^0 \rightarrow \Sigma^+ e^- \bar{\nu}_e$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
-1.7 ^{+2.1} _{-2.0} ±0.5	487	¹ ALAVI-HARATI01I	KTEV	p nucleus, 800 GeV

¹ ALAVI-HARATI 01I thus assumes that $g_2 = 0$ in calculating g_1/f_1 , above.

 $f_2(0)/f_1(0)$ FOR $\Xi^0 \rightarrow \Sigma^+ e^- \bar{\nu}_e$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
2.0±0.9 OUR AVERAGE				
2.0±1.3		BATLEY	13	NA48 p Be, 400 GeV
2.0±1.2±0.5	487	ALAVI-HARATI01I	KTEV	p nucleus, 800 GeV

 Ξ^0 REFERENCES

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ABLIKIM	23AP	PR D108 012006	M. Ablikim <i>et al.</i>	(BESIII Collab.)
ABLIKIM	23AU	PR D108 L031106	M. Ablikim <i>et al.</i>	(BESIII Collab.)
ABLIKIM	23B	PR D107 012002	M. Ablikim <i>et al.</i>	(BESIII Collab.)
PDG	22	PTEP 2022 083C01	R.L. Workman <i>et al.</i>	(PDG Collab.)
BATLEY	13	PL B720 105	J.R. Batley <i>et al.</i>	(CERN NA48/1 Collab.)
BATLEY	10B	PL B693 241	J.R. Batley <i>et al.</i>	(CERN NA48/1 Collab.)
BATLEY	07	PL B645 36	J.R. Batley <i>et al.</i>	(CERN NA48/1 Collab.)
BATLEY	07C	PL B650 1	J.R. Batley <i>et al.</i>	(CERN NA48 Collab.)
ABOUZAID	05	PRL 95 081801	E. Abouzaid <i>et al.</i>	(FNAL KTeV Collab.)
WHITE	05	PRL 94 101804	C.G. White <i>et al.</i>	(FNAL HyperCP Collab.)
LAI	04A	PL B584 251	A. Lai <i>et al.</i>	(CERN NA48 Collab.)
JAMES	02	PRL 89 169901 (errat.)	C. James <i>et al.</i>	(MINN, MICH, WISC, RUTG)
TEIGE	02	PRL 89 169902 (errat.)	S. Teige <i>et al.</i>	(RUTG, MICH, MINN)
ALAVI-HARATI	01C	PRL 86 3239	A. Alavi-Harati <i>et al.</i>	(FNAL KTeV Collab.)
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FANTI	00	EPJ C12 69	V. Fanti <i>et al.</i>	(CERN NA48 Collab.)
AFFOLDER	99	PRL 82 3751	A. Affolder <i>et al.</i>	(FNAL KTeV Collab.)
JAMES	90	PRL 64 843	C. James <i>et al.</i>	(MINN, MICH, WISC, RUTG)
TEIGE	89	PRL 63 2717	S. Teige <i>et al.</i>	(RUTG, MICH, MINN)
HANDLER	82	PR D25 639	R. Handler <i>et al.</i>	(WISC, MICH, MINN+)
COX	81	PRL 46 877	P.T. Cox <i>et al.</i>	(MICH, WISC, RUTG, MINN+)
BUNCE	79	PL B68 386	G.R.M. Bunce <i>et al.</i>	(BNL, MICH, RUTG+)
BUNCE	78	PR D18 633	G.R.M. Bunce <i>et al.</i>	(WISC, MICH, RUTG)
ZECH	77	NP B124 413	G. Zech <i>et al.</i>	(SIEG, CERN, DORT, HEIDH)
GEWENIGER	75	PL 57B 193	C. Geweniger <i>et al.</i>	(CERN, HEIDH)
BALTAY	74	PR D9 49	C. Baltay <i>et al.</i>	(COLU, BING) J
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MAYEUR	72	NP B47 333	C. Mayeur <i>et al.</i>	(BRUX, CERN, TUFTS, LOUC)
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WILQUET	72	PL 42B 372	G. Wilquet <i>et al.</i>	(BRUX, CERN, TUFTS+)
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PALMER	68	PL 26B 323	R.B. Palmer <i>et al.</i>	(BNL, SYRA)
BERGE	66	PR 147 945	J.P. Berge <i>et al.</i>	(LRL)

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PJERROU	65B	PRL 14 275	G.M. Pjerrou <i>et al.</i>	(UCLA)
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CARMONY	64B	PRL 12 482	D.D. Carmony <i>et al.</i>	(UCLA)
HUBBARD	64	PR 135 B183	J.R. Hubbard <i>et al.</i>	(LRL)
JAUNEAU	63	PL 4 49	L. Jauneau <i>et al.</i>	(EPOL, CERN, LOUC+)
Also		Siena Conf. 1 1	L. Jauneau <i>et al.</i>	(EPOL, CERN, LOUC+)